

A NEUROLOGICAL MODEL FOR SHELL PATTERN FORMATION

by

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An Honors Thesis  
Mathematical Sciences  
George Mason University

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A Neurological Model for Shell Pattern Formation

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## **Dedication**

I dedicate this thesis to all my brothers and sisters in arms fighting around the world. I have worked and trained beside many of you and all of you always remain in my heart. I also know the toll it takes on all of your loved ones and I hope they stay strong for you. I pray for peace and your safe return.

## Acknowledgments

There are numerous persons I would like to thank. To begin with, I would like to thank my parents for raising me to be a strong individual and for showing me the importance of education. They supported Karlys and myself while we pursued our undergraduate degrees. Without them, we would have never gotten to this point in our lives. I would like to thank my younger brother Zachary for giving me someone to fight with as a younger child, and wish him the best of luck with his first child which will be born in July. I would like to thank my wife, Karlys, for always encouraging me and supporting me in every endeavor I have pursued. I would also like to thank her for being brave enough to support me in my decision to leave the military to pursue my education in Mathematics.

I would like to thank the United States Airforce for showing me what my life would be like if I never strived to attain a higher position in life. I would like to thank my seventh grade Life Science Teacher from Francis C. Hammond Jr. High School Mrs. Bort. She was the first person to tell me I had potential ( what she actually said was, "you have a good brain, why don't you try using it for a change?"). I would like to think I took her advice. I would like to thank Mrs. Nancy Chovin of Antelope Valley College who recognized my potential in mathematics and encouraged me to follow it. I would like to thank all the wonderful Professors of Mathematics at George Mason University not only for the great education they bestowed upon me, but for also being very approachable and giving me great advice through the years. Specifically, I would like to thank Dr. Evelyn Sander for being my adviser and holding my hand through this process. Also, I would like to thank Dr. Sander, Dr. Goldin, and Dr. Wanner for

being such great teachers and believing in me enough to recommend me to graduate school. Finally, I would like to thank Christine Amaya and Pat Hendrix for making the Mathematics Department a welcome and comfortable place for all students.

I would also like to thank all the friends I have made over the years for their support and encouragement. Specifically, I would like to thank Keith Fox, Catherine Sausville, Samah Mahmoud, Sharif Thompson, Jason Youngblood, and Chris Vo for being such close friends for the last few years. I would like to give a special thanks to Robert Allen for encouraging me in writing my thesis and helping me edit it. Finally, I would like to thank my committee members Dr. Sander, Dr. Sauer, and Dr. Wanner for taking the time to help finish my thesis.

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# Abstract

## A NEUROLOGICAL MODEL FOR SHELL PATTERN FORMATION

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George Mason University, 2006

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The pattern formations of mollusk shells have been examined for millennia. We consider a neurological model for the central ganglion control pigmentation in shell pattern formation which is discrete in time and continuous in space. We will assume that the pattern is governed by the central ganglion and the amount of pigmentation to be secreted at a particular point in space and time is nonlocally coupled to all points in space at the previous time. This model is nonlinear. Therefore previous analyses are based on linearization and using eigenfunction expansions to approximate solutions. In this project we are studying the degree of error that comes from using these approximation techniques rather than the full nonlinear system. It is our hope that the results of this analysis will lead us to be able to isolate why certain patterns are selected over others.

# Chapter 1: Introduction

The brilliant patterns of sea shells have fascinated mankind for millennia. Unlike animal coat patterns, the patterns of sea shells formed by the molluscs living inside them do not appear to serve any defensive purpose. This is because molluscs spend almost their entire lives buried in the mud. So instead, scientists look for a neurological explanation to the formation of these patterns.

In this thesis, we will be using the neurological model for shell pattern formation originally proposed by Ermentrout, Campbell, and Oster[3]. We examine the different types of patterns which can be formed based on this model. We use a linearized version of this model to predict which types of patterns arise based on the eigenvalues. Also, we compare the linear solutions with the nonlinear solutions, and determine how long we can expect the linear solution to accurately reflect the nonlinear solution.

## 1.1 The Mollusc

In studying the formation of patterns on the shells of molluscs, it helps to understand some things about the creatures who create the shells. The class of mollusc's we will be concerning ourselves with is the class gastropoda. We should know a little gastropods form their shells. The first thing to note is that a "gastropod is a living, growing animal, and the opening of the shell is proportional to its body size. The shell grows with the body. New shell is laid down at the aperture, and reflects the

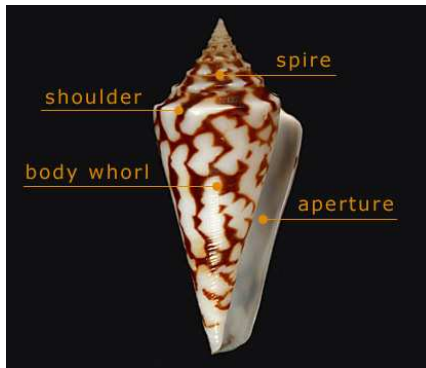


Figure 1.1: Basic Gastropod Shell[8]

size of the growing body”[7]. From Figure 1.1, we can see that the aperture of the shell is the opening of the shell. The shell of a gastropod takes the form of a spiral, so the growth of a shell starts from a central point and spirals outward. Now we will talk about some of anatomical features of gastropods.

This thesis is concerned with a neurological model for shell pattern formation, thus we should know a little about the nervous system of the gastropod. Figure 1.2 shows an image of the inner workings of a particular gastropod. In this model, we will be concerned with the central ganglia of the mollusc. The central ganglion is not listed in Figure 1.2, but we believe the central ganglion is just a combination of the all the various ganglions, such as the cerebral ganglion. Also from Figure 1.2, we can see the mantle of the shell. In the model we will be working with in this thesis, we will assume that the shell is formed at the mantel edge. Now that we have a basic understanding of gastropods, we can begin studying the model.

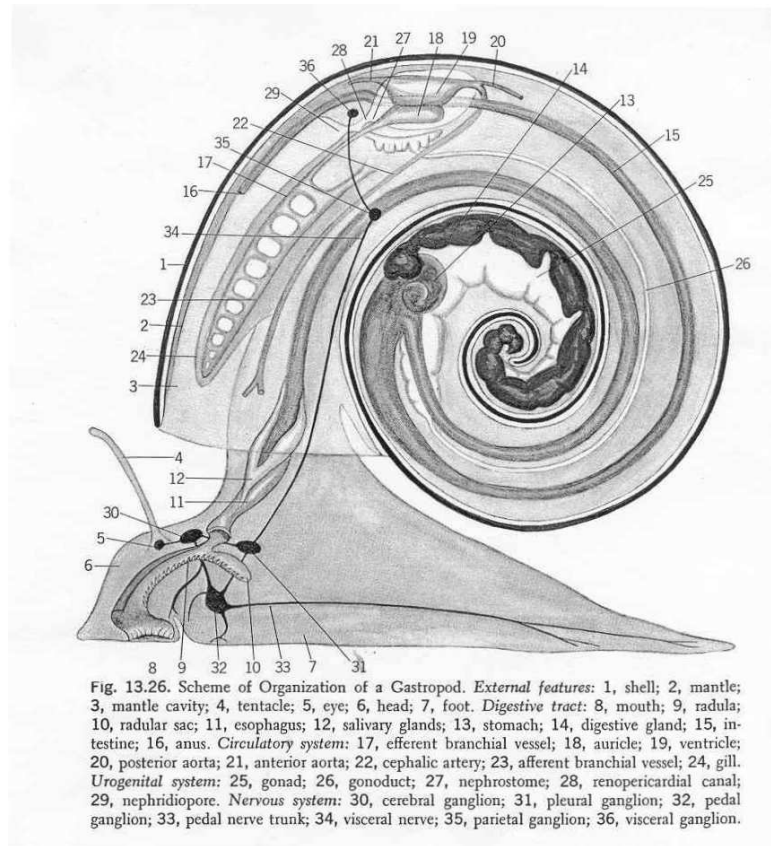


Figure 1.2: Gastropod Anatomy[7]

## 1.2 The Model

The model proposed by Ermentrout, Campbell, and Oster has the following assumptions[6]:

1. Cells at the mantle edge secrete material intermittently.
2. The secretion depends on (a) the neural stimulation,  $S$ , from surrounding regions of the mantle, and (b) the accumulation of an inhibitory substance,  $R$ , present in the secretory cell.



3. The net neural stimulation of the secretory cells consists of the difference between the excitatory and inhibitory inputs from the surrounding tissue.

From these assumptions, we need a functional  $S : \mathbb{R} \times C^0(\Omega) \rightarrow \mathbb{R}$  to determine the neural stimulation of the surrounding regions of the mantle. By  $C^0(\Omega)$ , mean the set of all continuous functions from  $\Omega$  to  $\mathbb{R}$ . We will also need a function  $R(t, x) : \mathbb{N} \cup \{0\} \times \mathbb{R} \rightarrow \mathbb{R}$  which gives the amount of inhibitory substance at time  $t$  and location  $x$ . The model also tells us that the amount of pigmentation,  $P(t, x) : \mathbb{N} \cup \{0\} \times \mathbb{R} \rightarrow \mathbb{R}$  laid down at a particular time  $t$  and space  $x$  is the difference in the neural stimulation and inhibitory substance. Based on this definition, we get the following recurrence relationship for  $P$ :

$$P(t + 1, x) = S(x, P(t, x)) - R(t, x). \quad (1.1)$$

We will also assume that  $R$  “depends linearly on the amount of pigment secreted during the previous period while at the same time degrading at a constant rate  $\delta$ ” [6]. The recurrence relationship for  $R$  is

$$R(t + 1, x) = \gamma P(t, x) + \delta R(t, x), \quad (1.2)$$

where  $0 < \gamma < 1$  is the rate of increase of pigmentation material, and  $0 < \delta < 1$  is the rate of decrease of inhibitory material.

Now we will define the neural stimulation functional  $S$ . Based on our assumptions,  $S$  is dependent on the surrounding regions of the mantle. We will further assume that the areas closer to the mantle will provide more stimulation than those which are further away. Thus, we need a filtering function which we will call  $w_j, j \in \{E, I\}$  where  $w_E$  is the excitatory filter, and  $w_I$  is the inhibitory filter. We will define

$w_j : \mathbb{R} \rightarrow \mathbb{R}$  in the following way:

$$\begin{aligned} w_j(x) &= 0 \text{ for } |x| > \sigma_j, j \in \{E, I\}, \\ w_j(x) &= q_j \left\{ 2^p - \left[ 1 - \cos \left( \frac{\pi x}{\sigma_j} \right) \right]^p \right\} \text{ for } |x| \leq \sigma_j, j \in \{E, I\}. \end{aligned} \tag{1.3}$$

Equation (1.3) was defined this way in [3]. However, a truncated exponential function would also work but would be discontinuous. The graph of this equation for a particular value of  $q_j$  and  $\sigma_j$  can be seen in Figure 1.3. The parameter  $q_j$  can be thought of as a scaling factor that when multiplied to the integral in  $w_j$  gives us  $\alpha_j$  as seen in equation (1.4). The parameter  $p$  can be thought of as controlling the sharpness of the cut-off.

$$\int_{\Omega} w_j(x) dx = \alpha_j, j \in \{E, I\}. \tag{1.4}$$

Therefore, we choose  $q_j$  based on the parameter value  $\alpha_j$ , which is the amplitude of the excitatory and inhibitory kernel functions. Also,  $\sigma_j$  is the measure of the range of the kernels. Using  $w_j$  as our filter, we now must apply it to  $P$ . Let us define  $E : \mathbb{R} \times C^0(\Omega) \rightarrow C^0(\Omega)$  to be the convolution of  $f$  with  $w_E$ ,

$$E(x, f) = \int_{\Omega} w_E(|x' - x|) f(x') dx' = (w_E * f)(x). \tag{1.5}$$

Let us define  $I : \mathbb{R} \times C^0(\Omega) \rightarrow C^0(\Omega)$

$$I(x, f) = \int_{\Omega} w_I(|x' - x|) f(x') dx' = (w_I * f)(x). \tag{1.6}$$

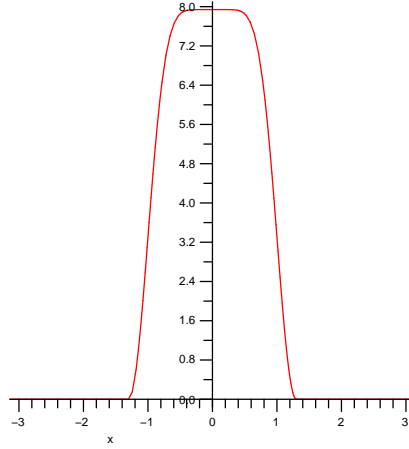


Figure 1.3: The  $w_j$  function for fixed  $q_j$  and  $\sigma_j$

Thus we are just performing a convolution on  $P$  and  $w_j$ . Now we can define the neural stimulation functional  $S : \mathbb{R} \times C^0(\Omega) \rightarrow \mathbb{R}$  as follows

$$S(x, f) = S_E(E(x, f)) - S_I(I(x, f)),$$

$$S_j : \mathbb{R} \rightarrow \mathbb{R}, \tag{1.7}$$

$$S_j(u) = \{1 + e^{-v_j(u - \theta_j)}\}^{-1}, j \in \{E, I\}.$$

We see that  $S_E(E(x, f))$  and  $S_I(I(x, f))$  are well defined for a fixed  $x$ . Clearly from the definition, the function  $S$  is nonlinear, making the model itself nonlinear. Thus, we will need to use linearization methods to perform analysis on the model. This will be done later in Chapter 3.

We can reduce equations (1.1) and (1.2) into one recurrence relation. To do this, we will begin by taking equation (1.1) and taking one step forward in time. So we

get,

$$P(t+2, x) = S(x, P(t+1, x)) - R(t+1, x). \quad (1.8)$$

Now take equation (1.2) and substitute it into equation (1.8). This give us,

$$P(t+2, x) = S(x, P(t+1, x)) - (\gamma P(t, x) + \delta R(t, x)). \quad (1.9)$$

Going back to equation (1.1), we can solve for  $R(t, x)$  which gives us  $R(t, x) = S(x, P(t, x)) - P(t+1, x)$ . Take this value for  $R(t, x)$  and plug that into equation (1.9) which gives us,

$$P(t+2, x) = S(x, P(t+1, x)) - (\gamma P(t, x) + \delta(S(x, P(t, x)) - P(t+1, x))). \quad (1.10)$$

Rearranging terms our recurrence relation becomes,

$$P(t+2, x) = S(x, P(t+1, x)) + \delta P(t+1, x) - \delta S(x, P(t, x)) - \gamma P(t, x). \quad (1.11)$$

Equation (1.11) give us a recurrence relation for our model which does not depend on  $R$ . From equation (1.11), we will generate our shell patterns. However, before we begin this process, we must discuss the numerical techniques used to generate these patterns and the accuracy of these methods.

## Chapter 2: Numerical Methods

The numerical calculations performed in this thesis use a combination of custom programs written in C++ , and the use of software packages such as Matlab and Maple. The C++ programs are used to calculate the actual patterns using pre-built libraries such as FFTW(Fastest Fourier Transform in the West). The Maple software package was used to for stability analysis of the nonlinear system by using linear approximations. Matlab was used to plot the patterns and aid in error analysis.

### 2.1 The Programs

A total of of three C++ programs were written for this thesis. The first program calculated the nonlinear pattern, the second program calculated linearized pattern, and final program made a comparison of the two patterns. Most of the programs are very similar; their differences will be mentioned when they are relevant.

#### 2.1.1 Homogeneous Equilibrium

For all the programs, we start by calculating a homogeneous equilibrium. First, a definition of a homogeneous equilibrium.

**Definition 2.1** (Homogeneous Equilibrium). An equilibrium is a point  $P(t, \Omega)$  such that

$$P(t_1, \Omega) = P(t_2, \Omega), \forall t_1, t_2 \in \mathbb{N} \cup \{0\},$$

so  $P(t, x)$  is unvarying with respect to  $t$ , so we can say that  $P(t, x) = P(x)$ . A homogeneous equilibrium is an equilibrium point  $P(x)$  such that

$$P(x) = P(y), \forall y \in \Omega.$$

So  $P(t, x)$  is unvarying with respect to  $t$  and  $x$ , so we can say that  $P(t, x) = P_0$  for some constant  $P_0 \in \Omega$ .

### Computing the Homogeneous Equilibrium

To compute the homogeneous equilibrium for the discrete time case, we start with our basic recurrence relation

$$P(t+2, x) = S(x, P(t+1, x)) + \delta P(t+1, x) - \delta S(x, P(t, x)) - \gamma P(t, x)$$

To find a homogeneous equilibrium, we have to apply the fact that is constant in space and time. Let us call our homogeneous equilibrium  $P_0$ . Therefore, we can rewrite our recurrence relation as

$$P_0 = S(x, P_0) + \delta P_0 - \delta S(x, P_0) - \gamma P_0. \quad (2.1)$$

Equation (2.1) is nonlinear, so a one-dimensional Newton's Method is used to solve for  $P_0 \in \mathbb{R}$ .

#### 2.1.2 Convolution

The model requires the use of the convolution of two functions in the computation of the patterns. We require the use of numerical techniques to compute the convolution, since in general, it does not have a closed form solution.

## Big Theta Notation

To understand the difficulty posed in computing the convolution of two functions, we will begin by defining what is meant by “Big-Theta Notation:”

**Definition 2.2** ( $\Theta$ -Notation[2]). For a given function  $g(n)$ , we denote by  $\Theta(g(n))$  the set of all functions

$$\Theta(g(n)) = \{f(n) : \exists c_1, c_2 \in \mathbb{R}, c_1, c_2 > 0, \text{ and } \exists n_0 \in \mathbb{N} \text{ such that}$$

$$0 \leq c_1 g(n) \leq f(n) \leq c_2 g(n), \forall n \geq n_0, n \in \mathbb{N}\}.$$

Let us say we have a function  $f : \mathbb{N} \rightarrow \mathbb{R}^+$  (where  $\mathbb{R}^+$  are the positive real numbers) such that  $f(n)$  represents the number of operations required to perform an algorithm on  $n$  data elements. For example, suppose we have a certain algorithm that will require  $f(n) = 3n^2 + n$  operations to sort an array of length  $n$ . It can be easily shown that  $f(n) \in \Theta(n^2)$  by choosing  $c_1 = 1$  and  $c_2 = 1000$ .

## The Discrete Fourier Transform

Direct computation of the convolution using the standard methods of numerical integration have a runtime complexity of  $\Theta(n^2)$ . Because of the large number of points required to approximate this infinite dimensional model, the convolution operation becomes the bottleneck of the program. Thus, in order to improve the efficiency of the computation, we need to find another method of computing the convolution. The following theorem provides us with an alternative.

**Theorem 2.1** (Convolution Theorem[4]). Let  $f, g$  be two functions with Fourier transforms, denoted  $\mathcal{F}(f), \mathcal{F}(g)$ , then the convolution of  $f, g$ , denoted  $f * g$  has the

following relation

$$f * g = \int_{\Omega} f(t')g(t - t')dt' = \mathcal{F}^{-1}(\mathcal{F}(f) \cdot \mathcal{F}(g)),$$

where  $\mathcal{F}(f) \cdot \mathcal{F}(g)$  is the pointwise product in Fourier space.

Thus we can perform the convolution in linear time in Fourier space. However, moving back and forth from Real space to Fourier space can bring about a new set of problems. The direct method of computing the Discrete Fourier Transform of a set of points has a runtime of  $\Theta(n^2)$ , where  $n$  is the number of points in the discretization of our domain  $[0, 2\pi]$ . However, there is a more elegant way of computing the Discrete Fourier Transform that will increase the runtime complexity significantly.

The algorithm, known as the Fast Fourier Transform, is capable of computing the Fourier transform of a set of  $n$  points with a runtime complexity of  $5n \log_2 n \in \Theta(n \log n)$ [5](Note, it can be shown that  $\Theta(n \log n) \cap \Theta(n^2) = \emptyset$ ). This improvement in speed with the Fast Fourier Transform algorithm comes with certain restrictions:

- all points must be equally spaced,
- values must be periodic,
- the number of points must be a power of two.

Note that if the points are not periodic, we can always extend it to be periodic. So, it would take  $2(5n \log_2 n) = 10n \log_2 n$  computations to convert the functions to Fourier space. The pointwise product then takes  $n$  computations, followed by an inverse Fourier Transform which adds an additional  $5n \log_2 n$  computations. Thus, the entire convolution operation using this method will take  $15n \log_2 n$  computations.



Assuming the direct method takes  $n^2$  operations, and choosing  $n = 2^{12}$ , we expect to see a speed up of over 23 times when compared to the direct method. Table 2.1 compares the run time of the convolution method using Fourier Transforms with the direct method.

Table 2.1: Comparison of Direct Convolution to Fourier Transform Convolution

Number of Points	Direct Method	Fourier Method
$2^{10}$	.2454s	.2718s
$2^{11}$	.5448s	.4414s
$2^{12}$	1.434s	.8314s
$2^{13}$	2.278s	1.4335s
$2^{14}$	6.3224s	1.4688s

### Fastest Fourier Transform in the West

The fact that we do not see a speed up of 23 times in Table 2.1 has to do with the initialization of the FFTW library, which performs check to find the optimal way to perform the Fourier transform. However, if we were to continue to increase the number of points, we would eventually see a speed up of at least 23 times.

The Fourier Transforms in this thesis were computed using a C library known as the “Fastest Fourier Transform in the West” (FFTW). The FFTW library was developed at MIT by Matteo Frigo and Steven G. Johnson. The FFTW library is a standard method of computing the Fourier Transform used for example by Matlab. However, we need to know the accuracy of this library in calculating the convolution. The issue of accuracy is because the functions  $w_E, w_I$  have a very small support in general. In fact, if we do not choose enough points, we could end up missing the support all together.

To test the accuracy of the FFTW compared, we compared the results to those obtained directly by using the trapezoid method of numerical integration. For a fixed set of parameters, we determined the  $\|\cdot\|_\infty$  of the differences for a set of common points.

Table 2.2: Error Comparison of Direct Convolution to Fourier Transform Convolution, (FT = Fourier transform, DM = direct method)

Method, Number of Points	FT, $2^{11}$	DM, $2^{11}$	FT, $2^{12}$	FT, $2^{13}$
FT, $2^{11}$	-	0.0018	$1.68 \cdot 10^{-5}$	$2.54 \cdot 10^{-5}$
DM, $2^{11}$	0.0018	-	0.0018	0.0018
FT, $2^{12}$	$1.68 \cdot 10^{-5}$	0.0018	-	$8.45 \cdot 10^{-6}$
FT, $2^{13}$	$2.54 \cdot 10^{-5}$	0.0018	$8.45 \cdot 10^{-6}$	-

In this way, we are measuring the convergence rate of the solution. It is clear that the Fourier transform is converging faster than the direct method. Thus, with the method of Fourier Transforms, not only do we get a time speed up, but we also get a more accurate computation of the convolution. Now that we have the method for computing the homogeneous equilibrium, along with a method for computing the convolution, we will describe the process for generating the shell patterns.

### 2.1.3 Pattern Generation Using C++

The pattern generation code was written in C++. In order to generate these patterns, we must approximate this infinite dimensional system with finitely many points. As would be expected, the more points we used, the more accurate the approximation. In this thesis, we have time range from time  $t = 0$  to time  $t = 1999$ . Since our domain is  $\Omega = [0, 2\pi]$ , we take  $2^{12} = 4096$  evenly spaced points on  $\Omega$ . So, 8192000

total points are used in generating each pattern. Now we can discuss the generation of the patterns.

In order to generate a pattern, we first start by computing the homogeneous equilibrium as described earlier. In some cases there are multiple homogeneous equilibria. Next, we take a nonhomogeneous random perturbation about the homogeneous equilibrium. This is done by adding a random amount to each coordinate in the domain. These values are generated randomly, and they are on the order of  $10^{-3}$ .

The first two time-steps,  $P(0, \Omega)$  and  $P(1, \Omega)$ , are generated this way. We need to do this because we are going to compute the pattern using Equation (1.11) which computes each new time step based on the previous two time-steps. These first two time-steps are the initial conditions for our system. We proceed iteratively to compute each new time-step as specified by Equation(1.11). If our homogeneous equilibrium is stable, and the perturbations stay within the basin of attraction, then the values will be attracted back to the homogeneous equilibrium. In this case, there is no pattern formation. However, if the homogeneous equilibrium is not stable, then we can see some very interesting patterns. These patterns will be discussed in detail in Chapter 3.

## 2.2 Analysis Using Maple

Maple is a program which was developed to perform symbolic computer based calculations. Maple uses a combination of symbolic manipulation techniques, along with numerical techniques to perform its computations. Maple also provides high level language constructs which can be used to create new Mathematical routines. In this thesis, we use Maple to compute the spectrum of the linearization of patterns

generated by a given set of parameters.

The nonlinear nature of the model makes predicting the behavior extremely difficult. Therefore, a way to simplify the analysis of this model is needed. This is done by linearizing the model and using the linear analysis to predict the behavior of the nonlinear model. Even with a linearized model, analysis can be difficult. Therefore, we used Maple to develop a program that when given a set of parameters, determines the eigenvalues of the model and the type of bifurcation for the nonlinear model. The Maple code can be seen in Appendix A.2.

## 2.3 Matlab

To end the numerical section of this thesis, we will discuss the other software program that was widely used. Matlab was used for many purposes; the most common uses were generating the patterns based on the data obtained from the C++ program, and aiding in error analysis and debugging. The code used to generate the patterns from the data produced by both the linear and nonlinear code, as well as graph of the difference between the linear and nonlinear code, can be seen in Appendix A.3.

## Chapter 3: Model Analysis

Thus far we introduced the model and showed the numerical methods that were used in computing the shell patterns. Now we will begin the discussion of the analysis of the model and results of that analysis.

### 3.1 Analysis of the Discrete Time Model

This discussion follows [6]. The discrete time model is a nonlinear system of equations which makes any direct analysis difficult. Thus, to perform analysis on the model we will linearize the model about a homogeneous equilibrium.

#### 3.1.1 Linearization of the Discrete Time Model

Let us take  $P_0$  to be a homogeneous equilibrium. We will linearize about  $P_0$  by writing

$$P(t, x) = P_0 + u(t, x), |u(t, \Omega)| \text{ small} . \quad (3.1)$$

Now we will substitute this into equation (1.11) we get

$$\begin{aligned} P_0 + u(t + 2, x) = & S(x, P_0 + u(t + 1, x)) - \gamma(P_0 + u(t, x)) - \\ & - \delta(S(x, P_0 + u(t, x)) - (P_0 + u(t + 1, x))). \end{aligned} \quad (3.2)$$

The only nonlinear contribution in equation (3.2) is the functional  $S$ , so we will now linearize  $S$ . To do this, we compute the Taylor series expansion about  $P_0$  for

$S(x, P_0 + u(t, x))$  and only keep the first order terms. Thus,

$$\begin{aligned}
 S(x, P_0 + u(t, x)) &\approx S(x, P_0) + DS(x, P_0)(u(t, x)) = \\
 &= S(x, P_0) + S'_E(E(x, P_0))DE(x, P_0)(u(t, x)) - \\
 &\quad - S'_I(I(x, P_0))DI(x, P_0)(u(t, x)).
 \end{aligned} \tag{3.3}$$

Since  $E(x, P(t, x)) = \int_{\Omega} w_E(|x' - x|)P(t, x')dx'$ , and since  $P_0$  is constant  $\forall x' \in \Omega$ , then

$$E(x, P_0) = \int_{\Omega} w_E(|x' - x|)P_0 dx' = P_0 \int_{\Omega} w_E(|x' - x|)dx' = \alpha_E P_0. \tag{3.4}$$

Likewise,  $I(x, P(t, x))$  for a fixed  $P_0$  we will get that

$$I(x, P_0) = \int_{\Omega} w_I(|x' - x|)P_0 dx' = P_0 \int_{\Omega} w_I(|x' - x|)dx' = \alpha_I P_0. \tag{3.5}$$

Now let us determine what  $DS(x, P_0)(u(t, x))$ .

$$\begin{aligned}
 DS(x, P_0)(u(t, x)) &= S'_E(E(x, P_0))DE(x, P_0)(u(t, x)) - \\
 &\quad - S'_I(I(x, P_0))DI(x, P_0)(u(t, x)).
 \end{aligned} \tag{3.6}$$

Evaluating  $DE(x, P_0)$  we get

$$\begin{aligned}
DE(x, P_0)(u(t, x)) &= \lim_{\epsilon \rightarrow 0} \frac{(w_E * (P_0 + \epsilon u(t, x)))(x) - (w_E * P_0)(x)}{\epsilon} = \\
&= \lim_{\epsilon \rightarrow 0} \frac{(w_E * P_0)(x) + \epsilon(w_E * u(t, x))(x) - (w_E * P_0)(x)}{\epsilon} = \\
&= \lim_{\epsilon \rightarrow 0} \frac{\epsilon(w_E * u(t, x))(x)}{\epsilon} = \\
&= (w_E * u(t, x))(x).
\end{aligned} \tag{3.7}$$

Similarly, we get  $DI(x, P_0)(u(t, x)) = (w_I * u(t, x))(x)$ . Finally, we get that

$$S(x, P_0 + u(t, x)) \approx S(x, P_0) + S'_E(\alpha_E P_0)(w_E * u(t, x))(x) - S'_I(\alpha_I P_0)(w_I * u(t, x))(x). \tag{3.8}$$

Let us note that Equation (3.8) is different than the cooresponding equation in both [6] and [3]. They both made a mistake differentiating the functional  $S$  and came up with  $S'_j(P_0)$  instead of  $S'_j(\alpha_j P_0)$  for  $j \in \{E, I\}$ .

To simplify notations, we will define a functional  $L_0$ , which will be the convolution operator, as follows

$$L_0(x, u(t, x)) = S'(\alpha_E P_0)(w_E * u(t, x))(x) - S'(\alpha_I P_0)(w_I * u(t, x))(x). \tag{3.9}$$

From (3.8) we see that  $S(x, P_0 + u(t, x)) \approx S(x, P_0) + L_0(x, u(t, x))$ . Now, substituting back into equation (3.3) we get

$$\begin{aligned}
P_0 + u(t + 2, x) &= S(x, P_0) + L_0(x, u(t + 1, x)) - \gamma(P_0 + u(t, x)) - \\
&\quad - \delta(S(x, P_0) + L_0(x, u(t, x)) - (P_0 + u(t + 1, x))).
\end{aligned} \tag{3.10}$$

Finally, we can subtract equation (3.3) from equation (3.10) to get the recurrence relation

$$u(t+2, x) - L_0(x, u(t+1, x)) - \delta u(t+1, x) + \delta L_0(x, u(t, x)) + \gamma u(t, x) = 0. \quad (3.11)$$

Now that we have the linearization of the model, we will discuss the analysis of the linear model.

### 3.1.2 Analysis of Linearized Model

To perform the analysis of the linear model, we will assume that  $u(t, \Omega) \propto \lambda^t e^{ikx}$ . Substituting this back into equation (3.11) we get

$$\lambda^{t+2} e^{ikx} - \lambda^{t+1} L_0(x, e^{ikx}) - \delta \lambda^{t+1} e^{ikx} + \delta \lambda^t L_0(x, e^{ikx}) + \gamma \lambda^t e^{ikx} = 0. \quad (3.12)$$

Now, let us see what happens with  $L_0(x, e^{ikx})$

$$\begin{aligned} L_0(x, e^{ikx}) &= S'_E(\alpha_E P_0) w_E * e^{ikx} - S'_I(\alpha_I P_0) w_I * e^{ikx} \\ &= S'_E(\alpha_E P_0) \int_{\Omega} w_E(|x' - x|) e^{ikx'} dx' \\ &\quad - S'_I(\alpha_I P_0) \int_{\Omega} w_I(|x' - x|) e^{ikx'} dx' \\ &= S'_E(\alpha_E P_0) \int_{\Omega} w_E(|x|) e^{ik(x-x')} dx' \\ &\quad - S'_I(\alpha_I P_0) \int_{\Omega} w_I(|x|) e^{ik(x-x')} dx' \\ &= e^{ikx} [S'_E W_E(k) - S'_I W_I(k)] \equiv e^{ikx} L^*(k), \end{aligned} \quad (3.13)$$

where we are defining  $W_E(k)$  and  $W_I(k)$  to be the Fourier transforms of  $w_E(x)$  and



$w_I(x)$  over the domain  $\Omega$  defined by

$$W_j(k) = \int_{\Omega} w_j(x) e^{-ikx} dx, \quad j = E, I. \quad (3.14)$$

Now we substitute equation (3.13) into equation (3.12) to get

$$\lambda^{t+2} e^{ikx} - \lambda^{t+1} e^{ikx} L^*(k) - \delta \lambda^{t+1} e^{ikx} + \delta \lambda^t e^{ikx} L^*(k) + \gamma \lambda^t e^{ikx} = 0. \quad (3.15)$$

Each term in equation (3.15) has a common factor of  $\lambda^t e^{ikx}$ , so by canceling out these terms we get

$$\begin{aligned} \lambda^2 - \lambda L^*(k) - \delta \lambda + \delta L^*(k) + \gamma &= 0, \\ \lambda^2 - \lambda(L^*(k) + \delta) + (\delta L^*(k) + \gamma) &= 0. \end{aligned} \quad (3.16)$$

Letting  $a(k) = -(L^*(k) + \delta)$ , and letting  $b(k) = \delta L^*(k) + \gamma$  we get finally get

$$\lambda^2 + a(k)\lambda + b(k) = 0. \quad (3.17)$$

Solving equation (3.17) using the quadratic equation we get that the eigenvalues are

$$\lambda(k) = \frac{-a(k) \pm \sqrt{a^2(k) - 4b(k)}}{2}. \quad (3.18)$$

Since the model is discrete in time, stability is attained if and only if  $|\lambda(k)| < 1, \forall k$ . Therefore, we have a stable homogeneous equilibrium if and only if all the complex eigenvalues are contained within the unit circle in the complex plane. Another way to look at the stability of an equilibrium is to consider the (a,b)-plane. We want  $|\lambda| < 1$ ,

thus considering what occurs in the (a,b)-plane we need

$$|\lambda| = \left| \frac{-a \pm \sqrt{a^2 - 4b}}{2} \right| < 1. \quad (3.19)$$

First let us consider what happens if we consider  $a^2 - 4b < 0$ . This would give us complex eigenvalues. Thus, we can write

$$\left| \frac{-a \pm \sqrt{a^2 - 4b}}{2} \right| = \left| \frac{-a \pm i\sqrt{4b - a^2}}{2} \right|. \quad (3.20)$$

Since  $-a + i\sqrt{4b - a^2}$  and  $-a - i\sqrt{4b - a^2}$  are complex conjugates of each other, they have the same modulus. Therefore, we can write equation (3.19) as

$$\sqrt{a^2 + 4b - a^2} < 2 \implies b < 1. \quad (3.21)$$

Now, let us consider what happens if we have real eigenvalues, thus  $a^2 - 4b > 0$ . Therefore, we get that

$$|-a \pm \sqrt{a^2 - 4b}| < 2. \quad (3.22)$$

We will start with  $|-a + \sqrt{a^2 - 4b}| < 2$ , this give us that  $-2 + a < \sqrt{a^2 - 4b} < 2 + a$ . From this, we know that  $a > -2$ . Since  $\sqrt{a^2 - 4b} < 2 + a$ , then we can solve for  $b$  to get

$$b > -a - 1 \text{ with } a > -2 \quad (3.23)$$

Now, let us take  $-2 + a < \sqrt{a^2 - 4b}$ . If we square both sides, and solve for  $b$ , we will get  $b < a - 1$  and  $b > a - 1$ . Therefore,  $-2 + a < \sqrt{a^2 - 4b}$  is not stable for any value of  $(a, b)$ .

Let us take  $|-a - \sqrt{a^2 - 4b}| < 2 \implies |a + \sqrt{a^2 - 4b}| < 2$ . Thus,  $-2 - a <$

$\sqrt{a^2 - 4b} < 2 - a$ . This implies that  $a < 2$ , so now taking  $\sqrt{a^2 - 4b} < 2 - a$  and solving for  $b$  we get,

$$b > a - 1 \text{ with } a < 2 \quad (3.24)$$

Finally, consider  $-2 - a < \sqrt{a^2 - 4b}$ . If we square both sides, and solve for  $b$ , we will get  $b < -a - 1$  and  $b > -a - 1$ . Therefore,  $-2 - a < \sqrt{a^2 - 4b}$  is not stable for any value of  $(a, b)$ .

Taking the intersections of equations (3.21), (3.23), and (3.24) we see that the region of stability occurs inside of a triangle in the  $(a, b)$ -plane. Figure (3.1) shows the bounded regions. As you can see from Figure 3.1, there is a parabolic dividing

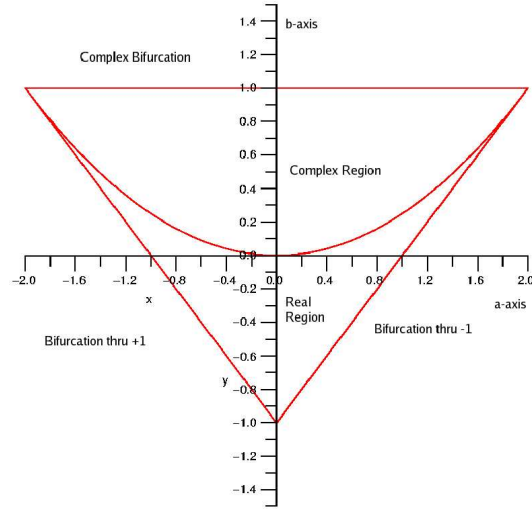


Figure 3.1: Bifurcation Triangle

line between the real region and the complex regions. This comes from the fact that when  $a^2 - 4b < 0$  we have complex solutions. Thus, when  $b > \frac{a^2}{4}$ , we have complex solutions, otherwise we have real solutions.

Now let us consider what occurs if we have a stable equilibrium. We can represent

all the eigenvalues of the system as a curve as a function of  $k$  in the  $(a, b)$ -plane, using equation (3.18). If the curve is entirely contained within the triangle shown in Figure 3.1, then we have stable homogeneous equilibrium. Now suppose we have a curve which leaves the stable region. The curve on the  $(a, b)$ -plane will be seen to be continuous. Therefore, the curve of the eigenvalues may intersect the triangle, but the actual eigenvalues only occur at discrete intervals. We know the eigenvalues occur at discrete time intervals because we are assuming periodic boundary conditions. So, in order to have an unstable equilibrium, we must have a  $k \in \mathbb{N}$  such that  $|\lambda(k)| > 1$  where  $\lambda : \mathbb{N} \rightarrow \mathbb{R}$  which will return the eigenvalue at  $k$ . Since there an  $|\lambda(k)| > 1$ , and we assumed that  $u(t, x) \propto \lambda^t e^{ikx}$ , and  $|\lambda^t e^{ikx}| = |\lambda^t|$ , and since  $|\lambda| > 1$ , then as  $t \rightarrow \infty$ ,  $\lambda^t \rightarrow \infty$  and so  $u(t, \Omega)$  must also go to  $\infty$ .

When the curve in the  $(a, b)$ -plane intersects the triangle in Figure 3.1, then we have a bifurcation at that point. A bifurcation “refers to significant changes in the set of fixed or periodic points or other sets of dynamic interest” [1]. The dynamic interest we are concerned with is when by changing the parameter we go from a having an attracting homogeneous equilibrium which creates no patterns, to having an unstable homogeneous equilibrium which will form a pattern. For this model, we will see three types of bifurcations. The first type is when there exists an eigenvalue  $\lambda > 1$ , we will call this a bifurcation thru +1. The second type occurs when we there exists an eigenvalue  $\lambda < -1$ , we will call this a bifurcation thru -1. The final type of bifurcation we are concerned with is when there exists an eigenvalue  $\lambda$  which is complex, such that  $|\lambda| > 1$ . Now we will examine how these bifurcations occur in Figure 3.1.

First we will consider when  $|\lambda| > 1$  where  $\lambda$  is complex occurs in the in Figure 3.1. Since  $\lambda$  is complex only when  $b > a^2/4$ , and we only have stable complex eigenvalues

when  $b < 1$ , then we will see  $|\lambda| > 1$  only when  $b > 1$ . So, a complex bifurcation will go through the upper part of the triangle in Figure 3.1

To see where a +1 bifurcation occurs in the  $(a, b)$ -plane, let us consider

$$\frac{-a(k) \pm \sqrt{a^2(k) - 4b(k)}}{2} > 1. \quad (3.25)$$

Solving Equation 3.25 for  $b$  we get two equations,  $b > -a - 1$  and  $b < a - 1$ , so this gives us

$$-a - 1 < a - 1 \implies a < 0. \quad (3.26)$$

Now we will show that -1 bifurcations occurs on the right side of the triangle in Figure 3.1, when  $b = a - 1$ . Plugging  $b = a - 1$  into equation (3.18) we get

$$\begin{aligned} \frac{-a \pm \sqrt{a^2 - 4(a - 1)}}{2} &= \frac{-a \pm \sqrt{(a - 2)^2}}{2} = \frac{-a \pm (a - 2)}{2} = \\ &= \begin{cases} -1 & \text{if } + \\ -a + 1 & \text{if } - \end{cases} \end{aligned} \quad (3.27)$$

As you can see, the eigenvalues on the line  $b = a - 1$  will always be -1, so a bifurcation through the right side of the triangle in Figure 3.1 will produce a bifurcation through -1.

## 3.2 +1 Bifurcation

Now we will begin to analyze what occurs when we have a bifurcation through  $\lambda = +1$ . From our previous analysis, we discovered that we will have a +1 bifurcation when the  $\lambda$  curve intersects the left side of the triangle in Figure 3.1. To study the types of patterns which form from these types of bifurcations, we will examine two parameter

sets which give us this type of bifurcation. These parameter can be seen in Table (3.1).

Table 3.1: Parameter Values for +1 Bifurcations

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	0.0	0.0	8.0	6.0	0.1	0.2	0.1	0.2	1.3	1.3
2	1.2	1.5	13.0	9.0	0.4	1.0	0.5	0.1	1.0	1.0

These parameter values were obtained using the Maple code by varying the parameters to come up with parameter values that worked. The bifurcations for these two parameter sets can be seen in Figure (3.2).

Now we will use the linearization to determine the types of patterns we will generate. We are assuming that our eigenvalues have the form  $\lambda^t e^{ikx}$ , since we are only concerned with positive real eigenvalues, then we can write this as  $\lambda^t \cos(kx)$ , where  $k$  is the dominate frequency. Thus, at any particular time step, we will see a sinusoidal pattern with frequency  $2\pi/k$ , with amplitude of  $\lambda^t$ . Since the frequency of the pattern does not change from one time period to the next, then we should expect to see striped pattern which is homogeneous in time. Now that we have some intuition about the types of patterns we should see, we will now examine the pattern formed by the first parameter set.

### 3.2.1 Parameter Set 1

In the original paper by Ermentrout *et al.*, a parameter set which has a +1 bifurcation was predicted to create “stationary spatial pattern of regularly spaced stripes”[3] by the linear model. For parameter set 1, we can see the linear solution in Figure 3.3. From the linear pattern, as well as the associated figure which shows time periods 40

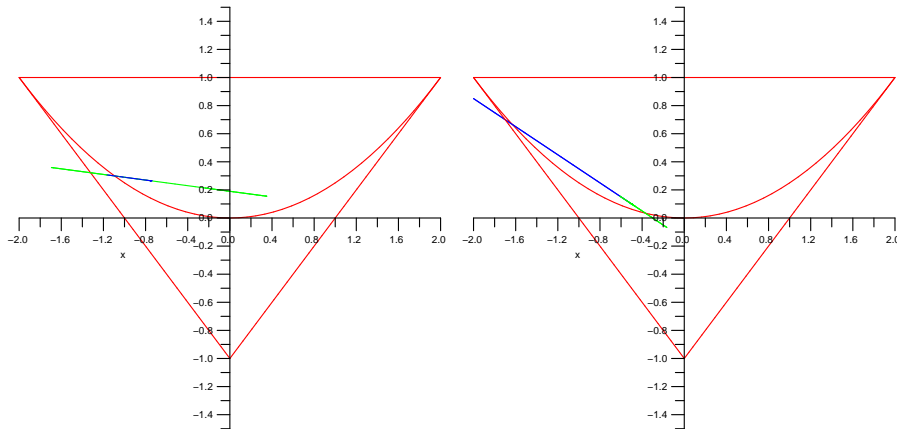


Figure 3.2:  $(a, b)$ -plane bifurcations for parameter set 1 and 2

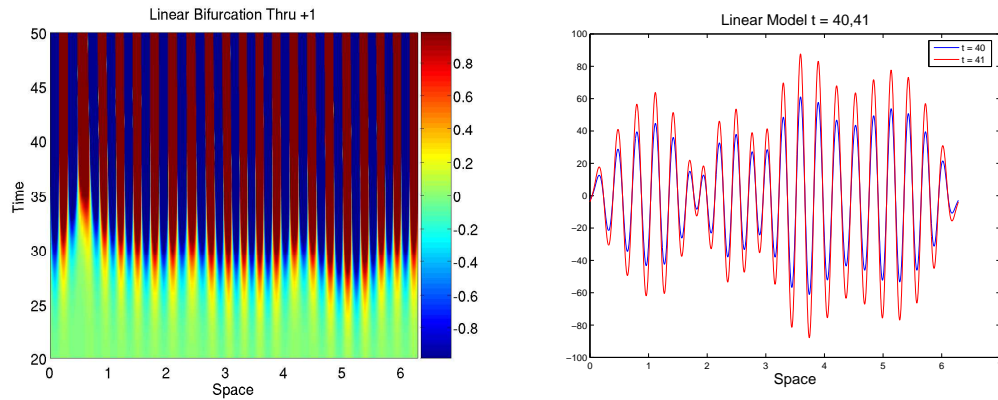


Figure 3.3: Linear Pattern for a +1 Bifurcation of parameter set 1, Linear Pattern at  $t = 40, 41$

and 41, we can see that there are 21 cycles, which gives us a wave length of  $2\pi/21$ . Since we have 21 wave cycles, we can expect to have 42 vertical stripes. Now if we examine Figure 3.4, we can see this is well within the range of possible values that the Maple code predicted.

Now let us consider the similarities and differences between the linear and non-linear patterns. As you can see, Figure 3.3 also has a stationary spatial pattern as

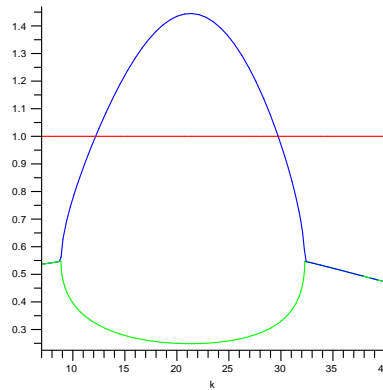
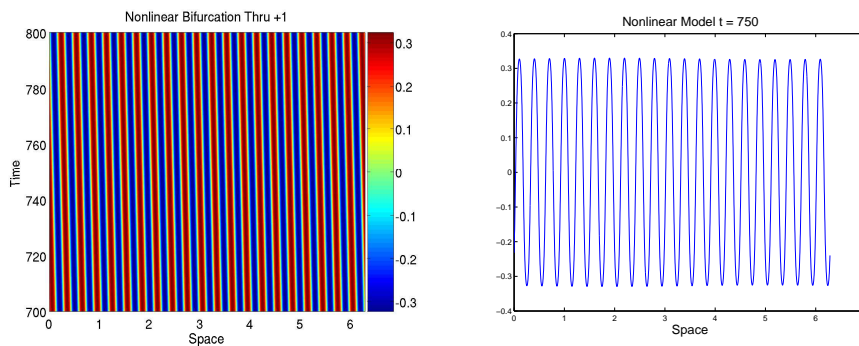


Figure 3.4: Eigenvalues for Parameter Set 1

expected. Now, if we examine the nonlinear pattern, we expect to see something similar. Figure 3.5 is the nonlinear pattern corresponding to the linear pattern shown

Figure 3.5: Nonlinear Pattern for a +1 Bifurcation of parameter set 1, Nonlinear Pattern at  $t = 750$ 

in Figure 3.3. Let us examine the similarities of this patterns. First, we can clearly see that both the linear and nonlinear patterns seem to have approximately stationary spatial patterns of regularly spaced stripes. Also, both the linear and nonlinear pattern has the same number of cycle, 21, and therefore the same number of vertical stripes, 42. However, there are some differences. Notice that the nonlinear pattern



appears to be moving slightly to the left as time increases. This is a nonlinear behavior that does not seem to appear in the linear pattern. This may be the result of some asymmetry that is inherent in the nonlinear model, which is not representable in the linear model. The other major difference between the linear and nonlinear patterns is that while the nonlinear solution is taken on some form of pattern, the linear pattern is diverging off to infinity. Therefore, we can only expect the linear pattern to give us a good approximation of the nonlinear pattern for a finite amount of time. Now, we want to be able to measure how long the linear model will approximate the nonlinear model for a given parameter set. We will do this by taking the  $\infty$ -norm of the differences of the linear and nonlinear model for each time  $t$ . We will say that the linear model approximates the nonlinear model as long as there difference is less than 1. Now let us look at Figure 3.6. As you can see from Figure 3.6, the linear

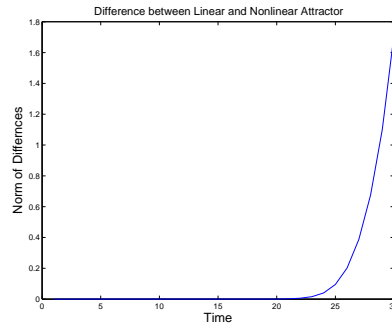


Figure 3.6: +1 Bifurcation Difference between Linear and Nonlinear

model appears to approximate the nonlinear model for about 29 time steps. We can now ask, how a slight change in the parameters affects the time period that the linear model approximates the nonlinear model. We will attempt a partial answer to this question later on. For now, let us look at the second parameter set.

### 3.2.2 Parameter Set 2

We have already seen from Figure 3.1 that the second parameter set also has a bifurcation through the second quadrant. Despite this, we will see a very different pattern from the first parameter set. First we will take a look at the linear pattern. It can be seen in Figure 3.7. Once again, we see the horizontal striped pattern, but

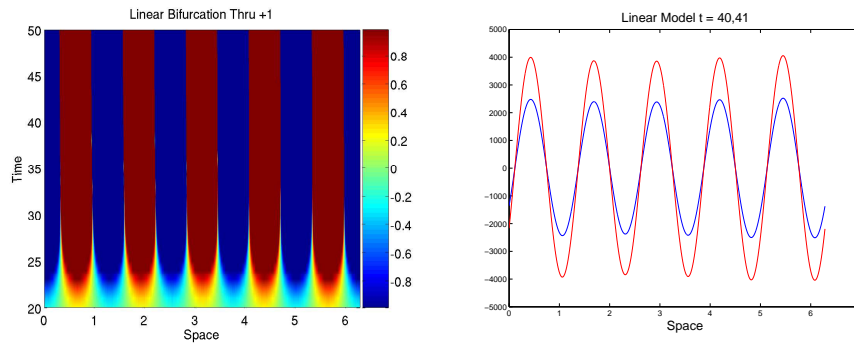


Figure 3.7: Linear +1 Bifurcation for Parameter Set 2, Linear Pattern at  $t = 40, 41$

this time there are only 10 stripes, significantly less than the 42 we saw in Figure 3.3. The reason for this is the same as the reason in the previous problem set, the dominate eigenvalues. Figure 3.8 shows us the positive eigenvalues. Now let us see

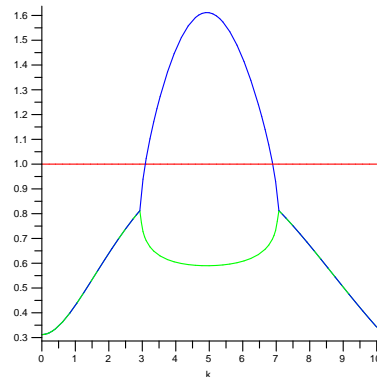


Figure 3.8: Eigenvalues for Parameter Set 2

what happens in this case with the nonlinear pattern. That pattern can be seen in Figure 3.9. Once again, we see the similarities between the linear(Figure 3.7) and

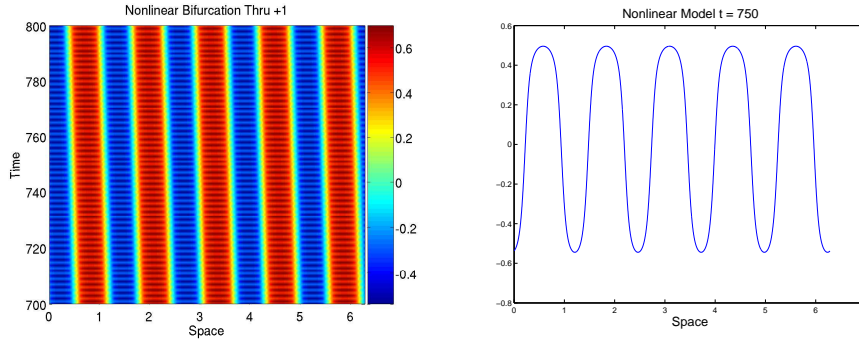


Figure 3.9: Nonlinear +1 Bifurcation for Parameter Set 2, Nonlinear Pattern at  $t = 750$

nonlinear(Figure 3.9) patterns. They both has 10 distinct stripes and they both take have vertical stripe patterns. However, again we see some differences

As we saw in Figure 3.5, Figure 3.9 also moves to the left as time increases. However, this time we also see that there appears to be some spatial variation in the vertical stripes of the nonlinear pattern. We do not see this at all in the linear pattern, meaning this again is nonlinear behavior that cannot be duplicated by the linear pattern.

To finish this section, we would once again like to see how long the linear pattern can approximate the nonlinear pattern. We can see this in Figure 3.10. This time we see that the linear pattern approximates the the nonlinear pattern for about 24 time iterates.

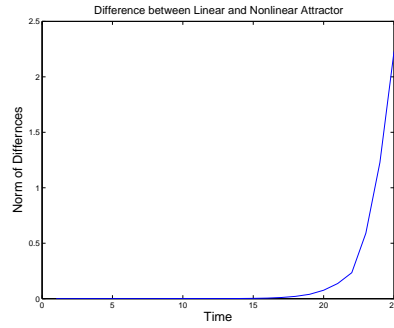


Figure 3.10: +1 Bifurcation Difference between Linear and Nonlinear for Parameter Set 2

### 3.2.3 Variation of Initial Conditions

One last topic we would like to explore before moving onto the next type of bifurcation is what sort of differences, if any, do we see when we vary the initial conditions. It turns out that if we vary the initial conditions of parameter set 2, we can end up with a different number of vertical stripes. This can be seen in Figure 3.11.

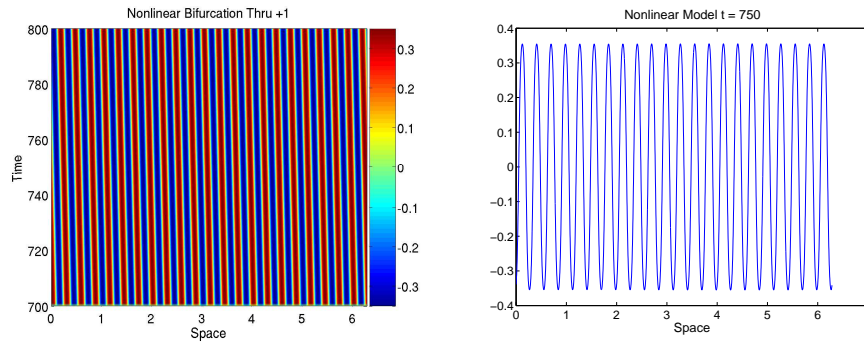


Figure 3.11: Nonlinear +1 Bifurcation of Parameter Set 1 with a Different Number of Vertical Stripes

In Figure 3.4 we see  $\lambda(k) > 1$  for both  $k = 21$  and  $k = 22$ . Since they are both unstable eigenvalues, it is possible for an eigenvalue which is not the largest eigenvalue to dominate for a period of time depending on the initial conditions. Note, that this

effect occurs in both the linear and nonlinear models, since it is only a result of the eigenvalues. Now we will discuss -1 bifurcations.

### 3.3 $-1$ Bifurcation

We know from before that  $-1$  bifurcations occurs when the  $\lambda$  curve intersects the triangle in Figure 3.1 on the right side of the triangle. Once again, to study these types of bifurcations we will choose two parameter sets which have -1 bifurcations. These parameter sets can be seen in Table 3.2.

Table 3.2: Parameter Values for -1 Bifurcations

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	2.2	0.0	8.0	6.0	0.1	.2	0.3	0.1	1.5	1.5
2	1.0	2.0	15.0	13.0	1.0	0.1	0.2	0.1	0.5	0.5

Since we are looking at eigenvalues  $\lambda < -1$ , we can write

$$\lambda^t e^{ikx} = (-1)^t |\lambda|^t \cos(kx),$$

since we are only concerned with the real values of  $\lambda$ . Therefore, we should expect to see a period two oscillation in time due to the  $(-1)^t$  value. Also, just like +1 bifurcations, we expect sinusoidal behavior in space because of the  $\cos(kx)$ . Now that we have some intuition of patterns should emerge, let us analyze the first parameter set.

#### 3.3.1 Parameter Set 1

First we will begin by looking at the bifurcation diagram in the  $(a, b)$ -plane. As you can see by Figure 3.12, the  $\lambda$  curve intersects the right side of the triangle.

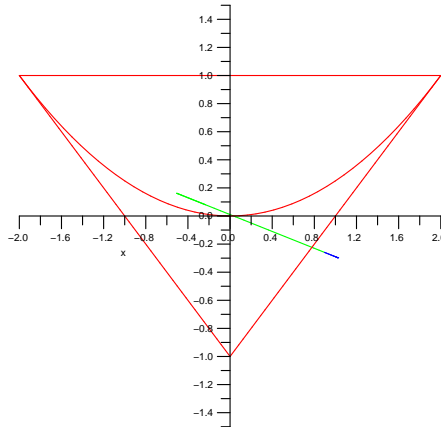


Figure 3.12:  $(a, b)$ -plane Bifurcations for Parameter Set 1

Since this is a -1 bifurcation, we know we have a sinusoidal behavior in space. We expect the wavelength of these patterns to be  $\frac{2\pi}{k_M}$ , where  $|\lambda(k_M)|$  is maximal. However, Murray tells that the wavelengths vary “with wavenumbers  $k$  in the range bounded by  $k_1$  and  $k_2$  with a maximum  $\lambda$ ,  $\lambda_M = \lambda(k_M)$ ” [6]. Here  $k_1$  and  $k_2 \in \mathbb{N} \cup \{0\}$ , where  $k_1$  is the smallest value in  $\mathbb{N} \cup \{0\}$  such that  $|\lambda(k_1)| > 1$  and  $k_2$  is the largest value in  $\mathbb{N} \cup \{0\}$  such that  $|\lambda(k_2)| > 1$ . So, in the long term behavior of the system, we would expect to have wavelength of length  $\frac{2\pi}{k_M}$ , however, at any particular time, we could see wavelength varying from  $k_1 < k < k_2$  where  $\lambda(k_1), \lambda(k_2) < -1$ . For the first parameter set, we have a bulk bifurcation, meaning that  $k_M = 0$ , as predicted by the Maple Code, see Figure 3.13.

Therefore, we would expect the long term behavior of the wave length to be “infinity,” which can be seen by the linear pattern shown in Figure 3.14. Figure 3.14 shows us that in this case, our pattern follows the behavior of  $k_M$ , however, we shall see that this not always the case.

Now, let us discuss how the linear pattern compares with the nonlinear pattern.

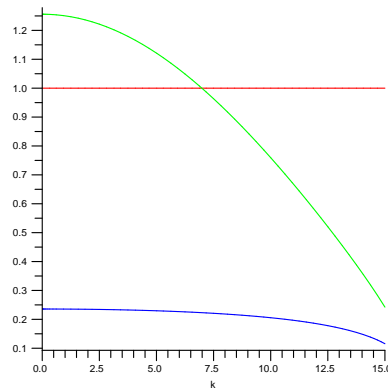


Figure 3.13: The eigenvalues for Parameter Set 1

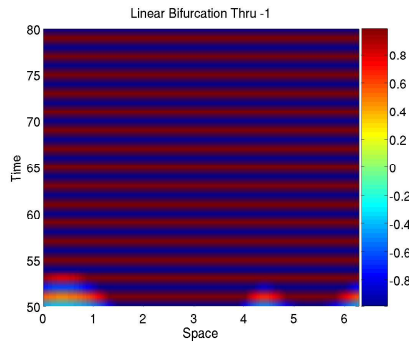


Figure 3.14: Linear -1 Bifurcation for Parameter Set 1

The nonlinear pattern can be seen in Figure 3.15

Finally, we can see how long the linear pattern approximates the nonlinear pattern by examining Figure 3.16

As you can see, the linear model for this parameter set approximates the nonlinear one for about 47 time iterates. Now we will examine the second parameter set.

### 3.3.2 Parameter Set 2

As before, we will begin our discussion by examining the bifurcation diagram for parameter set 2 in the  $(a, b)$ -plane. Now, we want to look at the range of the eigenvalues

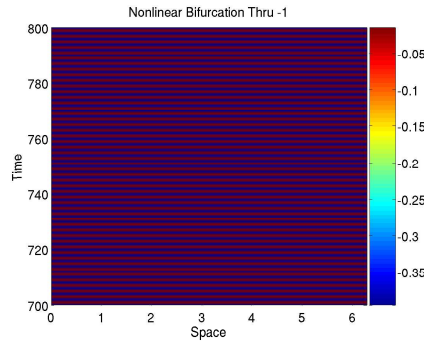


Figure 3.15: Nonlinear -1 Bifurcation for Parameter Set 1

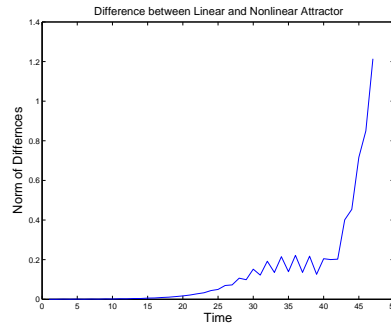


Figure 3.16: Difference Between the Linear and Nonlinear Patterns for Parameter Set 1

of this parameter set to see what kind of pattern to expect. Figure 3.18 will show us the range of the eigenvalues  $|\lambda| > 1$ . So, we can see from Figure 3.18 that we can expect our  $k$  value to be  $4 < k < 20$  with  $k_M \approx 6$ . Now let us look at Figure 3.19 to see what the linear pattern actually looks like. From Figure 3.19, we see that there are 12 vertical stripes, and by examining the linear pattern at  $t = 15, 16$  we see that there are exactly 6 cycles in that graph. This tell us that we have a wavelength of  $2\pi/6$  which was predicted by Murray. Now let us examine the nonlinear pattern and compare it to the linear pattern. We can see that the nonlinear pattern follows exactly from the linear pattern. We see that the nonlinear pattern is period 2 in time



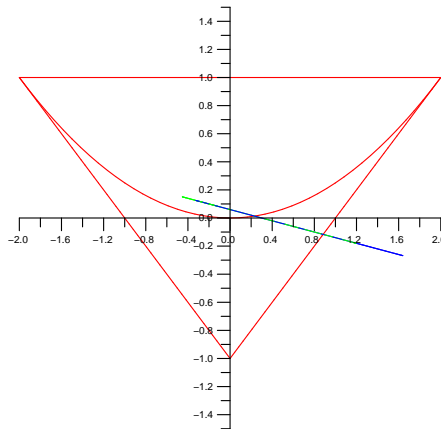


Figure 3.17:  $(a, b)$ -plane Bifurcations for Parameter Set 2

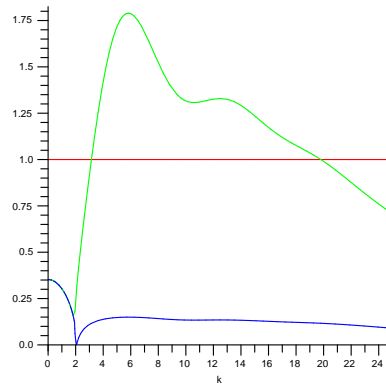


Figure 3.18: The eigenvalues for Parameter Set 2

and has 6 cycles in the graph. Now, let us consider the length of time that the linear model approximates the nonlinear model for this parameter set. Figure 3.21 shows that we can expect the linear pattern to approximate the nonlinear pattern for about the first 20 time iterates. Now we shall examine how the patterns can change based on a change in initial conditions, but leaving the parameters unchanged.

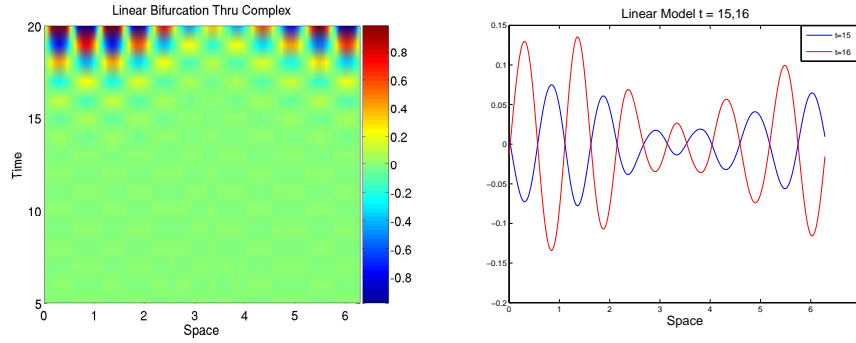


Figure 3.19: Linear -1 Bifurcation for Parameter Set 2, Linear pattern at  $t = 15, 16$

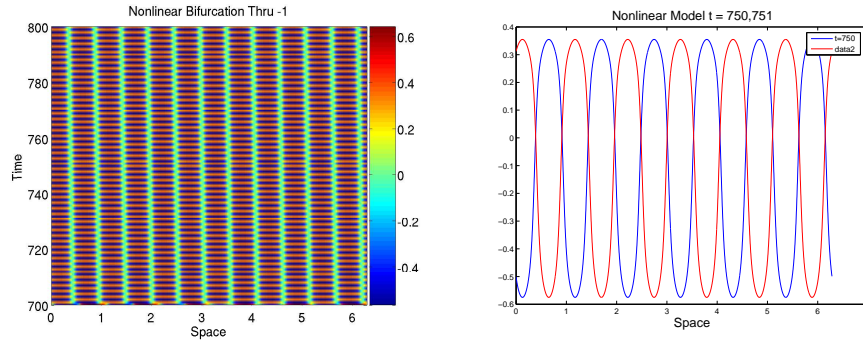


Figure 3.20: Nonlinear -1 Bifurcation for Parameter Set 2, Nonlinear pattern at  $t = 750, 751$

### 3.3.3 Variation of Initial Conditions

If we vary the initial conditions of parameter set 1, we can see a different pattern than the one shown in Figure 3.14. We know from Figure 3.13 that it is possible to have  $k$  values ranging from 0 to about 30. Now let us examine Figure 3.22. From this figure we can see that we have one cycle, implying that in this case we have  $k = 1$ . Therefore, just as in the +1 bifurcations, we can see a different eigenvalue from the most dominate one control how the pattern is formed for certain initial conditions.

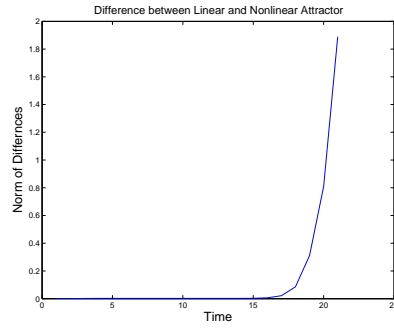


Figure 3.21: Difference Between the Linear and Nonlinear Pattern for Parameter Set 2

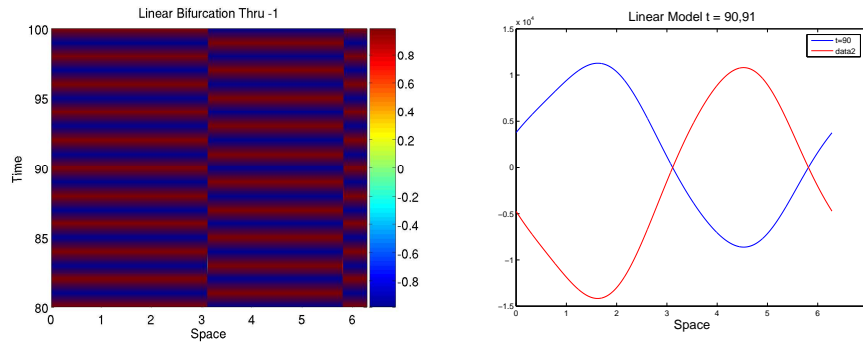


Figure 3.22: Linear -1 Bifurcation for Parameter Set 1 with a Different Pattern

### 3.4 Complex Bifurcation

Complex bifurcations occur when the  $\lambda$  curve intersects the top part of the triangle in Figure 3.1. To understand the patterns these types of bifurcations create, let us take a look at the parameter values in Table 3.3.

Table 3.3: Parameter Values for Complex Bifurcations

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	2.1	3.0	10.0	9.0	0.1	0.2	0.7	0.8	0.7	0.7
2	0.2	0.3	20.0	7.0	1.3	0.7	0.4	0.905	1.5	1.5

Since we have a complex bifurcation, then we have a  $\lambda^t = r^t e^{i\theta t}$ ,  $\theta \neq 0, \theta \neq \pi$ .

Since  $|\lambda| > 1$ , then  $r > 1$ , so we see exponential growth. However, since this is a complex bifurcation,  $\arg(\lambda) > 0$ . This tells us that the pattern travels on the unit circle by the amount  $\arg(\lambda)$ . Also, we still have the spatial variations as we had in +1 and -1 bifurcations. Because of this, we expect to see a spatio-temporal pattern [3]. Now, let us take a look at parameter set 1.

### 3.4.1 Parameter Set 1

First, we will look at the bifurcation graph for this parameter set. It can be seen in Figure 3.23. Now we will calculate the spatial variations we should expect from this

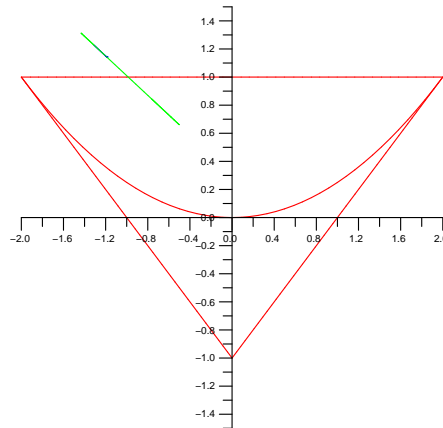


Figure 3.23:  $(a, b)$ -plane Bifurcations for Parameter Set 1

parameter set. From figure 3.24 we see that  $k_1 = 0$  and  $k_2 \approx 36$  with  $k_M \approx 22$ . Thus, we could see wavelengths anywhere from length “ $\infty$ ” to  $2\pi/36$ . However, since we have a complex bifurcation, we must also be concerned with the  $\arg(\lambda)$ . The  $\arg(\lambda)$  will determine the periodicity we see in time. We use the Maple code to compute the  $\arg(\lambda)$  and use this to compute the expected periodicity in time of this particular pattern. The range of periodicity with respect to the  $k$  values can be seen in Figure

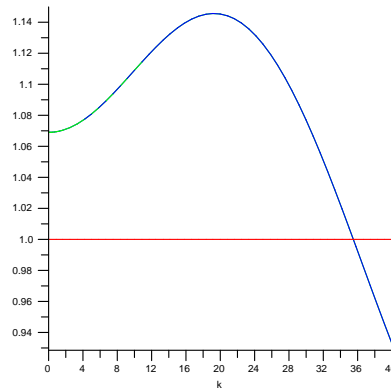
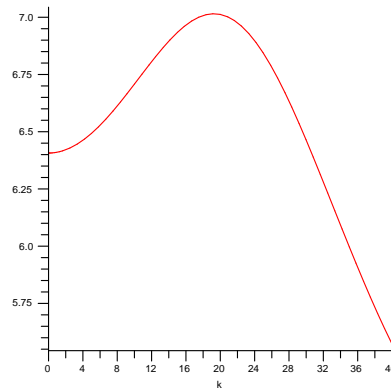


Figure 3.24: Eigenvalues for Parameter Set 1

3.25. We can expect the periodicity in time of parameter set 1 to a multiple  $n \in \mathbb{N}$

Figure 3.25: The  $2\pi/\arg(\lambda)$  for Parameter Set 1

of a number between 6 and 7, since these are the values of  $\arg(\lambda)$  when  $|\lambda| > 1$ . To determine the periodicity of our parameter set in time, we will take our nonlinear data and plot the  $t$  vs.  $t + 1$  for a fixed  $x$ . This gives us the periodicity in time because we expect our pattern to be periodic with period  $M$ . Thus, if we graph  $t$  vs.  $t + 1$  for a fixed  $x$ , then we will see at most a multiple of  $M$  points on the graph. The multiplicity  $n$  will be determined by the number of points crossed by each line

segment  $+1$  in the  $t$  vs.  $t + 1$  graph. Figure 3.26 shows us the  $t$  vs.  $t + 1$  plot for parameter set 1. From Figure 3.26, we see that no line segment crosses any points,

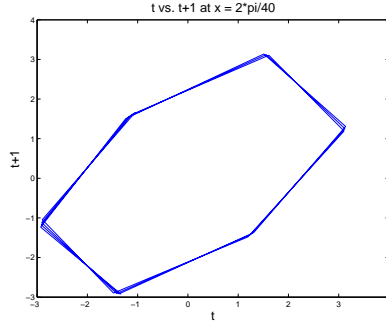


Figure 3.26: Periodicity of Parameter Set 1

thus the multiplicity of our parameter set is  $n = 1$ . Now, counting the number nodes in the graph, we see that the periodicity of this parameter set in time is slightly more than 6, which is what we would expect.

Let us see if this parameter set behaves as expected. Figure 3.27 show the linear pattern generated by parameter set 1, along with the 3 values  $u(t, \Omega)$  for  $t = 550, 551$ , and 552. We can see that there are 19 cycles in the linear pattern, cooresponding

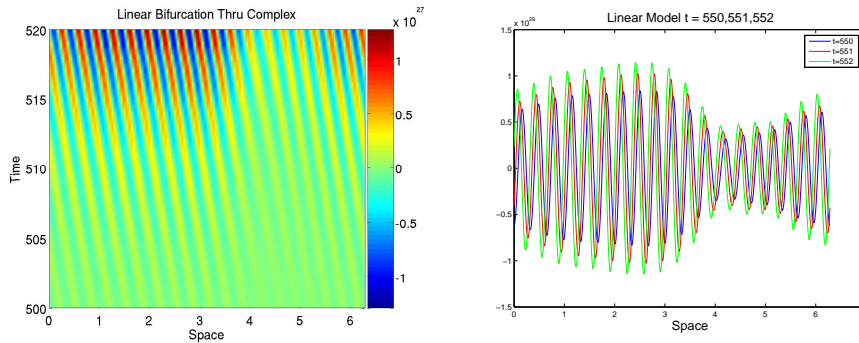


Figure 3.27: Linear Complex Bifurcation for Parameter Set 1, Linear Pattern at  $t = 550, 551, 552$

to a wavelength of  $2\pi/19$  which is in the predicted range. Also, we see that Figure

3.27 varies both spatially and temporally as expected. Now, let us examine the nonlinear pattern and see if it has the behavior we expect. Figure 3.28 has a spatio-

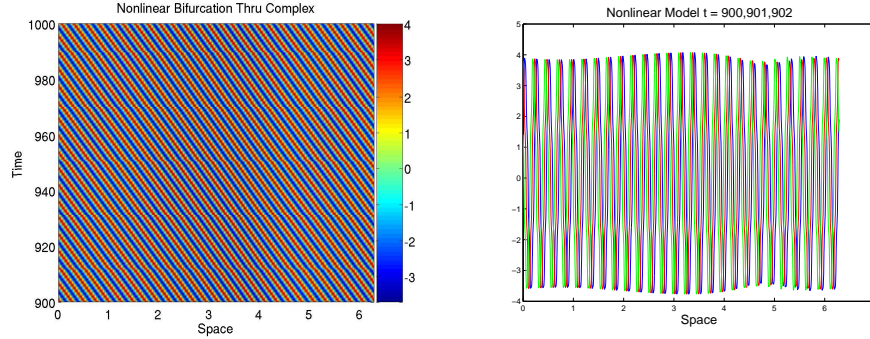


Figure 3.28: Nonlinear Complex Bifurcation for Parameter Set 1, Nonlinear Pattern at  $t = 900, 901, 902$

temporal pattern as expected. We can see that there are 24 cycles, corresponding to wavelengths of  $2\pi/24$ , which differs from the linear pattern. However, this wavelength is still well within the expected range of values. However, other than the difference in wavelengths, the linear and nonlinear patterns are identical.

Now we will look how long the linear pattern approximates the nonlinear pattern. Let us take a look at Figure 3.29. Thus, this linear pattern approximates the nonlinear

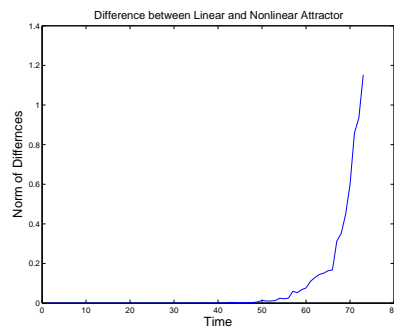


Figure 3.29: Difference Between the Linear and Nonlinear Patterns for Parameter Set

pattern for about 70 time iterates in this parameter set. Now, let us look at some interesting patterns that emerge from this parameter set.

Though this pattern seems simple, it actually has some interesting phenomena associated with it. In fact, in some of the earlier time periods, the pattern is actually much more complicated. See Figure 3.30 to view both the linear and nonlinear versions. As you can see, they both have this “V” shape, so it is not an oddity that

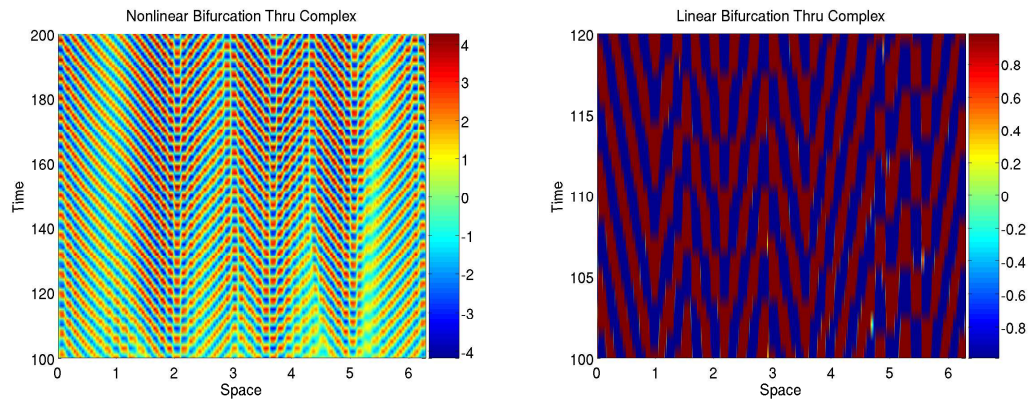


Figure 3.30: The “V” Shape for both the Nonlinear and Linear Models for Parameter Set 1

is only captured in the nonlinear case. These strange patterns can be credited to transience, and they are dampened out in the long run. Now let us take a look at the second parameter set.

### 3.4.2 Parameter Set 2

Let us begin again by first examining the bifurcation diagram. We can see the bifurcation diagram for the second parameter set in Figure 3.31. We can see that the  $\lambda$  curve intersects and crosses the upper part of the triangle, indicating that we have a complex bifurcation for parameter set 2. We will now use the Maple code to see



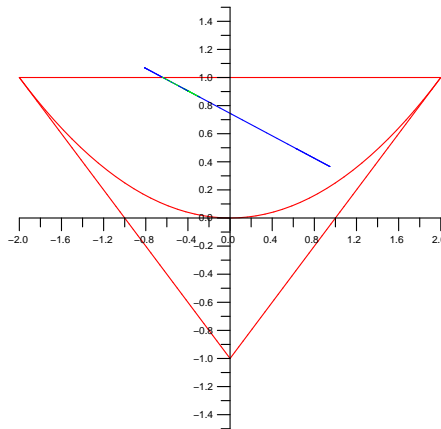


Figure 3.31:  $(a, b)$ -plane Bifurcation Diagram for Parameter Set 2

what type of pattern we should expect. Figure 3.32 we can see that  $k_1 \approx 7$ ,  $k_2 \approx 10$ ,

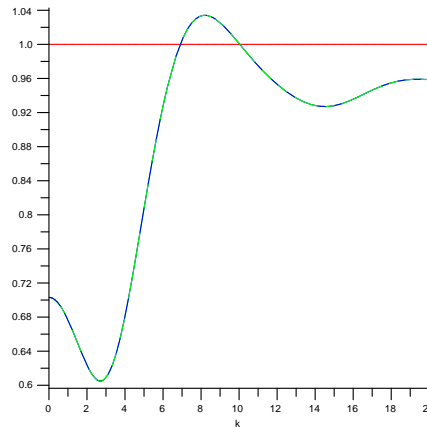


Figure 3.32: Eigenvalues for Parameter Set 2

and  $k_M \approx 9$ . So we should expect between 7 to 10 cycles in space. Since this is a complex bifurcation, we also want to use the  $\arg(\lambda)$  to give us some more intuition of what the pattern should look like. Figure 3.33 shows us that we should expect the periodicity of parameter set 2 to be a multiple of a number between 4.7 and 5.8. Now, let us take a look at the plot of  $t$  vs.  $t + 1$  which can be seen in Figure 3.34. We can see that this graph is approximately period 14. However, each edge crosses 2

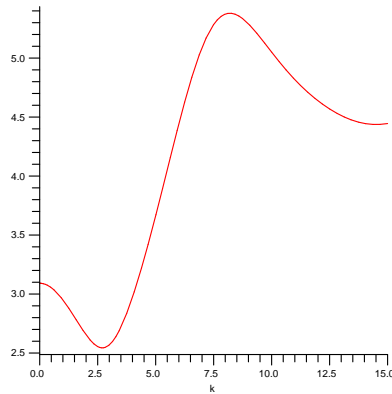


Figure 3.33: The  $2\pi/\arg(\lambda)$  for Parameter Set 2

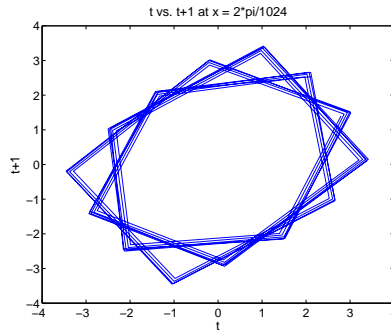


Figure 3.34: Periodicity of Parameter Set 2

points, thus the multiplicity of this graph is  $n = 3$ . Since  $4.7 \cdot 3 = 14.1 \approx 14$ , then we can see that the Maple code correctly predicted the periodicity of parameter set 2.

Now, we will normally we would examine the linear pattern, but there is a problem with the linear code for this parameter set. We will explain more about this later on. For now, let us examine the nonlinear pattern. Figure 3.35 shows us the nonlinear pattern along with the graph for fixed time periods 750, 751, 752, and 753. From the graph of the fixed time periods, we can see that there are 8 cycles in the graph, cooresponding to a wavelength of  $2\pi/8$  which is in the expected range.

To understand the problem we are having with generating the linear pattern for

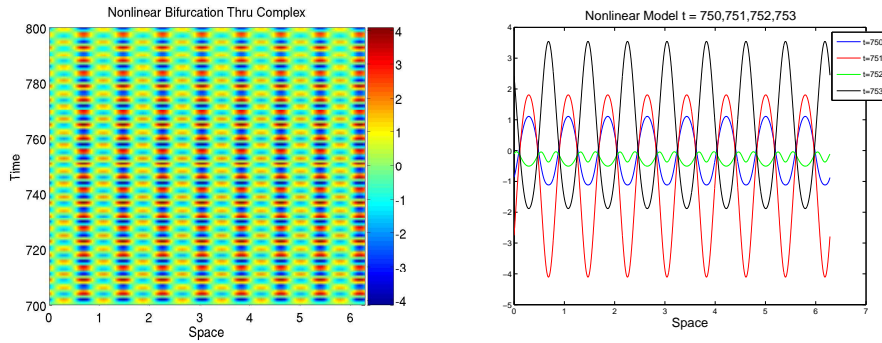


Figure 3.35: Nonlinear Complex Bifurcation of Parameter Set 2, Nonlinear Pattern at  $t = 750, 751, 752, 753$

this parameter set, let us first see how long the linear code is approximating the nonlinear code. From Figure 3.36 we can see linear and nonlinear patterns diverge

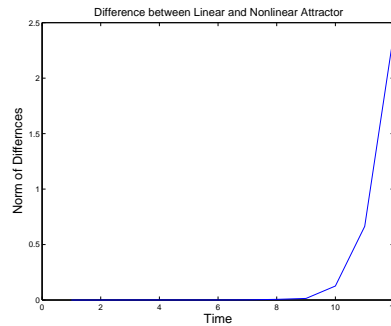


Figure 3.36: Difference Between Linear and Nonlinear Patterns for Parameter Set 2

after only 11 time iterates. The problem we are having with the linear code in this case is that we are not seeing any pattern formation. This could be caused by an error in the linear code, though this seems unlikely since the same linear code is used to produce all the other linear patterns. Another possibility is that the linear and nonlinear code are diverging so quickly that this somehow prevents the pattern from forming.

Now, we would like to see what will happen if we fix the parameters and vary the

initial conditions.

### 3.4.3 Variation of Initial Conditions

We are going to show that for a particular parameter set, a complex bifurcation can have different periodicity in time. To see how this occurs, let us look at Figure 3.37

Table 3.4: Parameter Values for Complex Bifurcations with Period 5 and 19

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	0.2	0.3	20.0	7.0	1.3	0.7	0.4	0.87	1.5	1.5

which show us the eigenvalues for this parameter set as well as the  $2\pi/\arg(\lambda)$ . We

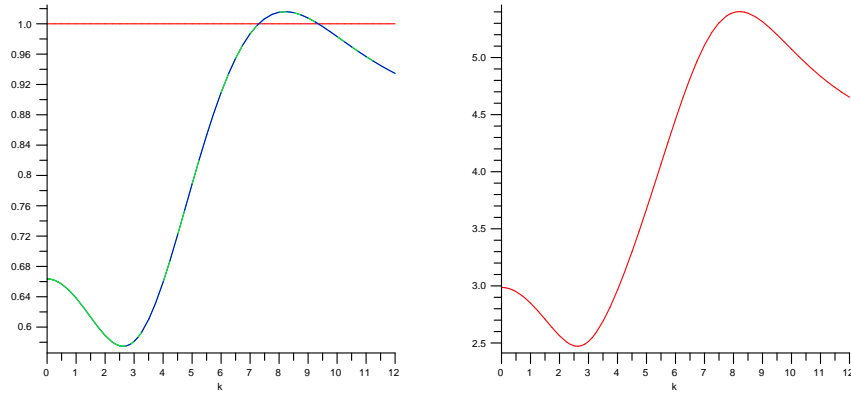


Figure 3.37: Eigenvalues and  $2\pi/\arg(\lambda)$  for the Parameter Set in Table 3.4

can see that the only  $k$  values that gives us  $|\lambda(k)| > 1$  are  $k = 7, 8, 9$ . Also, from the graph of  $2\pi/\arg(\lambda)$  we see that we should expect the periodicity of this parameter set to be approximately a multiple of a number between 4.8 to 5.8. For the first set of initial conditions, we can see the periodicity based on these initial conditions in Figure 3.38. We can see that the multiplicity is  $n = 1$  and we see a period 5 pattern as expected. Now, let us take a look at Figure 3.39 to see the periodicity of this

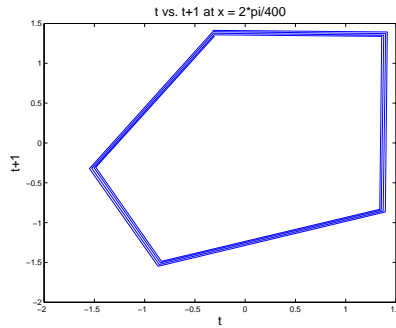


Figure 3.38: Initial Conditions which are Period 5 in Time.

parameter set for the second set of initial conditions. We finish this discussion by

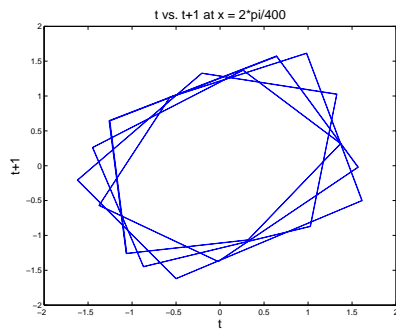


Figure 3.39: Initial Conditions which are Period 19 in Time.

We can see that the periodicity of this parameter set for the second set of initial conditions to be approximately period 19. We can see that each edge pass through 3 nodes, giving us a multiplicity of  $n = 4$ . This is consistent with the Maple code

$$\text{since } 4.8 \cdot 4 = 19.32 \approx 19.$$

examining the nonlinear patterns of this parameter set for the cooresponding initial conditions. These patterns are seen in Figure 3.40. Now, we will consider what happens when we have multiple bifurcations.

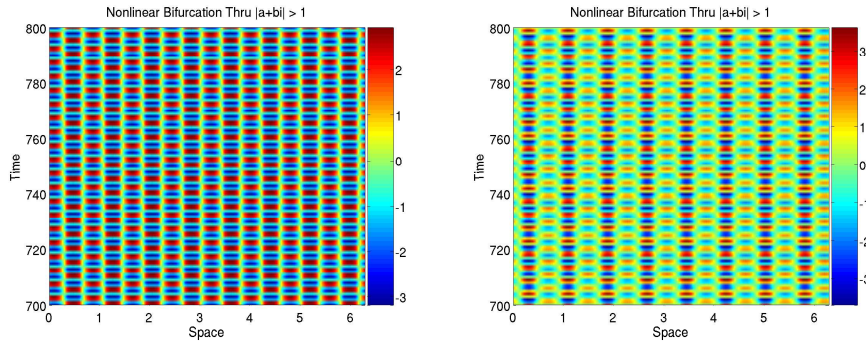


Figure 3.40: Patterns generated for  $\gamma = .87$  for a pattern which is period 5 in and period 19 in time.

## 3.5 Multiple Bifurcations

A parameter set which has  $\lambda$  curve in the  $(a, b) - plane$  which intersects the triangle in Figure 3.1 on multiple sides is said to have multiple bifurcations. We will not be doing an in-depth analysis of these types of bifurcations. Instead, we will show some of the interesting patterns these types of bifurcations creates. The first type of bifurcation we will see is a bifurcation through complex and -1.

### 3.5.1 Bifurcations thru Complex and -1

The parameter set that produced the pattern that we will be examining can be seen in Table 3.5. The nonlinear pattern for this parameter set can be seen in Figure 3.41

Table 3.5: Parameter Values for Bifurcation which are both Complex and -1

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	0.2	0.3	19.0	10.0	1.0	0.2	0.2	0.9	1.0	1.0

Now, we will be examining a pattern created from a +1 and -1 bifurcations.

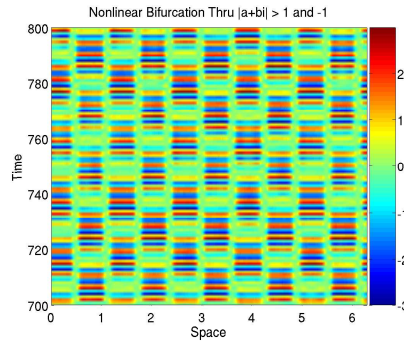


Figure 3.41: Nonlinear Bifurcation thru Complex and -1

### 3.5.2 Bifurcations thru +1 and -1

The parameter set used to generate the pattern we will be viewing can be seen in Table 3.6. The nonlinear pattern for this parameter set can be seen in Figure 3.42.

Table 3.6: Parameter Values for Bifurcation which are both +1 and -1

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	0.1	0.0	6.0	8.0	0.1	0.2	0.1	0.0	3.0	3.0

We conclude this thesis with a discuss of the rate of divergence of the linear and

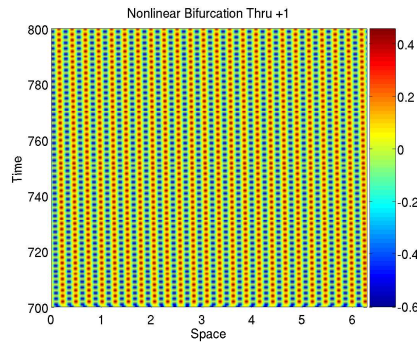


Figure 3.42: Nonlinear Bifurcation thru +1 and -1

nonlinear model.

### 3.6 Analysis of the Rate of Divergence

We will now consider the rate of divergence of the linear model with respect to the nonlinear model. To do this, we will fix all parameters except for a single parameter which will allow us to have an unstable solution bifurcate from a stable solution to a unstable solution. Table 3.7 shows the parameter set we will be using. This

Table 3.7: Parameter Values for Bifurcation which are both +1 and -1

Set Number	$\theta_E$	$\theta_I$	$\alpha_E$	$\alpha_I$	$\sigma_E$	$\sigma_I$	$\delta$	$\gamma$	$\nu_E$	$\nu_I$
1	2.1	3.0	10.0	9.0	0.1	0.2	0.7	variable	0.7	0.7

parameter set is the same as parameter set 1 in Table 3.3, except we are varying  $\gamma$ . The bifurcation remains a complex bifurcation as we vary  $\gamma$ . Figures 3.43 and 3.44 show us how the  $\lambda$  curve moves through the bifurcation as we vary  $\gamma$ . The plot

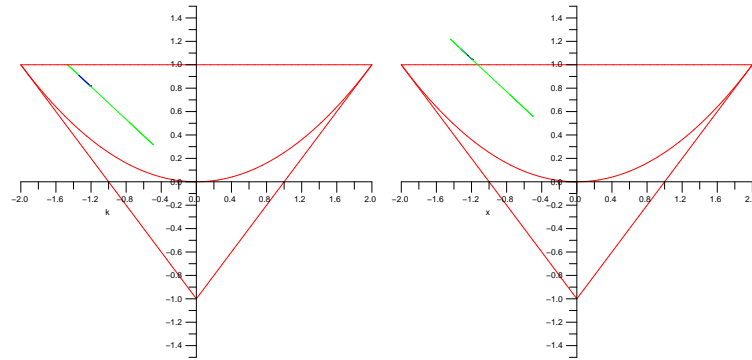


Figure 3.43: Complex Bifurcations for  $\gamma = .463$  and  $\gamma = .7$

the  $\|\cdot\|_\infty$  of the differences of the linear and nonlinear model with respect to  $\gamma$  can be seen in Figure 3.45. We see an exponential decay in the amount of time that the linear and nonlinear models approximate each other. This is because as the  $\lambda$  curves moves further through the complex bifurcation, the dominant eigenvalue, call it  $\lambda_M$ , increases in modulus. Thus, the rate at which the linear model diverges from the



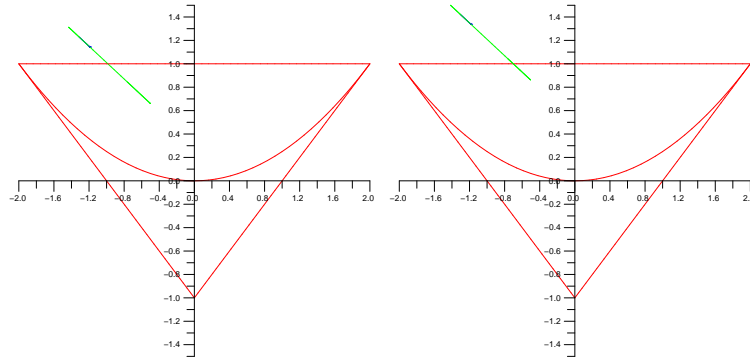


Figure 3.44: Complex Bifurcation for  $\gamma = .8$  and  $\gamma = .999$

nonlinear model increases exponentially as  $\gamma$  increases. In fact, this holds true for any type of bifurcation since the further the  $\lambda$  curve moves through the bifurcation, the larger the modulus of the maximum eigenvalues become.

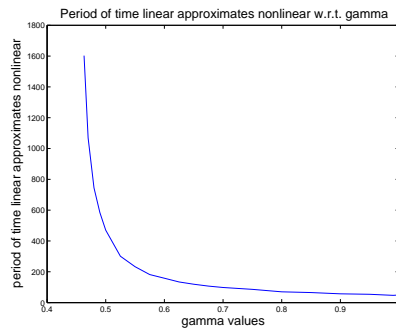


Figure 3.45: How  $\| \cdot \|_{\infty}$  of the Difference of the Linear and Nonlinear Model varies with respect to  $\gamma$

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## Appendix A: Source Code

### A.1 C++ Code

```

/*****

*   Copyright (C) 2005 by Tyler M. White                               *
*   twhite5@gmu.edu                                                    *
*                                                                 *
*   This program is free software; you can redistribute it and/or modify *
*   it under the terms of the GNU General Public License as published by *
*   the Free Software Foundation; either version 2 of the License, or   *
*   (at your option) any later version.                                *
*                                                                 *
*   This program is distributed in the hope that it will be useful,      *
*   but WITHOUT ANY WARRANTY; without even the implied warranty of      *
*   MERCHANTABILITY or FITNESS FOR A PARTICULAR PURPOSE.  See the      *
*   GNU General Public License for more details.                       *
*                                                                 *
*   You should have received a copy of the GNU General Public License    *
*   along with this program; if not, write to the                      *
*   Free Software Foundation, Inc.,                                       *
*   59 Temple Place - Suite 330, Boston, MA 02111-1307, USA.           *
*                                                                 *
*****/

/*Description:  This program was written by Tyler White to be used in his
/*              Honors Thesis in Mathematics.  This program implements an

```

```

/*          an algorithm used to produce shell patterns.  The algorithm
/*          this model was based on was original proposed by Ermentrout,
/*          Campbell, and Oster in 1986.  This code not only generates the
/*          the nonlinear patterns, but also computes the linear patterns
/*          based on the linearization of the nonlinear model and also
/*          computes the infinity norm of the differences between the
/*          linear and nonlinear patterns.*/

```

```

#ifdef HAVE_CONFIG_H

```

```

#include <config.h>

```

```

#endif

```

```

#include <iostream>      //for basic output to console

```

```

#include <math.h>        //for the basic C++ math libraries

```

```

#include <fstream>       //for the file stream class to write to files

```

```

#include <vector>        //for the Standard Template Library vector container

```

```

#include <fftw3.h>       //for the fast fourier transform library

```

```

#include <algorithm>     //for the Standard Template Library algorithms

```

```

#include <time.h>        //to get the system time

```

```

#include <stdlib.h>      //basic library functions such as rand()

```

```

#include <cstdlib>

```

```

using namespace std;

```

```

#define PI 3.14159265358979323846 //store highly accurate value for pi.

```

```

#define pwr 12           //used to set the size of the domain which will be 2^pwr

```

```

#define p 4

```

```

/*****
/*This section is the functions used by main to interface with the rest*/
/*of the program                                     */
void getInput();

//this function is used to get the input from the user

/*          */
void patternGen();

//this function is the interface between main and the computation part
//of the program

/*          */
void initializationVariables();

//this function initializes all the variables

void cleanUp();

//this function cleans up the variables at the end of the program

*****/

/*****
/*these functions are part of the original formulations of the model    */
double wE( double x );

//This function evaluates wE for all values of x in the domain

double wI( double x );

//This function evaluates wI for all values of x in the domain

double SE( double x );

//This function evaluates SE for all values of Pt(x) of x in the domain

double SI( double x );

```

```

//This function evaluates SI for all values of Pt(x) of x in the domain

double S( double x_E, double x_I );

//This function calls SE and SI and takes their difference and returns that
/*****/

/*****/

/*Functions used in determining the homogenous equilibrium solution */
/*                                                                    */
double equationNewton(double x);
//Part of newtons method, basically the f(x) of newtons method      */
/*                                                                    */
double equationNewtonPrime(double x);
//Part of newtons method, basically the fprime(x) of newtons method */
/*                                                                    */
double newtonMethodP0(double x);
//xt_p_1 = xt - f(x)/fprime(x)                                        */
/*                                                                    */
void computeEquilibrium();
/*uses fixed point iteration and newtons method to come up with a    */
/*solution.                                                            */
/*These are the variables used for calculating the equilibrium solution*/
double equilibrium = -0.3938;
double precision;
double accuracy = 8;

/*****/

/*****/

```

```

/*these functions are speifically used in the linearization */
/*
calculations*/

double SprimeE( double x );

//this function determines the derivative of SE for a given value of x

double SprimeI( double x );

//this function determines the derivative of SI for a given value of x

void computeL0();

//this function commputes the convolution of ut and wE and wI and uses

//this to compute L0

/*****/

/*****/

/*these are the auxillery fuctions which perform simple calculations */
/*help make the code more modular */

void arrayCpy( double a[], double b[] );

//Function that copies the values of array b into array a

void computeConvolution( fftw_plan&, fftw_complex[],

                        fftw_complex[], fftw_complex[] );

//This is used for computing the convolution operator

//It assumes that the variables have already been transformed into

//Fourier Space.

double infinity_norm(double*, double*);

//This uses the infinity norm to calculate the distance between two vectors

//in  $R^n$ , the first parameter is the Pt values and the second parameter is

//the Ut value.

```

```

double rel_norm(double*, double*, double);

/*****

/*****

/*These are the main parameters for this program      */
double theta_E = .2;
double theta_I = 0.3;
double alpha_E = 20;
double alpha_I = 7;
double sigma_E = 1.3;
double sigma_I = .7;
double gamma1 = .85;
double delta = 0.4;
double nu_E = 1.5;
double nu_I = 1.5;

double q_E = alpha_E*4/(93*sigma_E);    //parameter which depends on alpha,
                                         //p, and sigma
double q_I = alpha_I*4/(93*sigma_I);    //parameter which depends on alpha,
                                         //p, and sigma

const int N = 1 << pwr;                //evalutes 2^pwr using the bitshift operator
int max_time = 2000;                   //defines the length of the time interval, so
                                         //time goes from 0 to max_time

/*****

/*****

/*These are the fourier transform plans used to execute */
/*the fourier tranforms */

```



```

fftw_plan fft_wE;           //for the fourier transform of wE(x)
fftw_plan fft_Pt;           //for the fourier transform of Pt(x)
fftw_plan fft_wI;           //for the fourier transform of wI(x)
fftw_plan fft_Ut;
fftw_plan ifft_SwE;         //for the fourier transform of the conv of
                             //SE[Et(x)]
fftw_plan ifft_SwI;         //for the fourier transform of the conv of
                             //SI[It(x)]
fftw_plan ifft_SwE1;        //for the fourier transform of the conv of
                             //SE[Et(x)]
fftw_plan ifft_SwI1;        //for the fourier transform of the conv of
                             //SI[It(x)]

/*****

/*****

/*These are variables used in the fourier transforms for input and output*/

double *in_wE;              //input for the fft of wE(x)
double *in_wI;              //input for the fft of wI(x)
double *in_Ut;
fftw_complex *out_Ut;
fftw_complex *out_wE;       //output for the fft of wE(x)
fftw_complex *out_wI;       //output for the fft of wI(x)
double *in_Pt;              //input for the fft of Pt(x)
fftw_complex *out_Pt;       //output for the fft of Pt(x)
fftw_complex *in_SwE;        //input for the conv of Pt(x) and wE(x)
double *out_SwE;            //output for the conv of Pt(x) and wE(x)
fftw_complex *in_SwI;        //input for the conv of Pt(x) and wI(x)

```

```

double *out_SwI;          //output for the conv of Pt(x) and wI(x)
fftw_complex *in_SwE1;    //input for the conv of Pt(x) and wE(x)
double *out_SwE1;        //output for the conv of Pt(x) and wE(x)
fftw_complex *in_SwI1;    //input for the conv of Pt(x) and wI(x)
double *out_SwI1;        //output for the conv of Pt(x) and wI(x)

/*****/

/*****/

//these variables are used in the actual computations for Pt, Ut, and L0t

double *Pt;               //the pigmentation from 2 iterations ago
double *Pt_p_1;           //the pigmentation from 1 iteration ago
double *Pt_p_2;           //the currently calculated pigmentation
double *Ut;
double *Ut_p_1;
double *Ut_p_2;
double *L0;
double *L0t;
double *L0t_p_1;
double *L0t_p_2;
double x_coor[N];
double temp_SwE[N];
double temp_SwE1[N];
double temp_SwI1[N];
double temp_SwI[N];

/*****/

double diff = 2*PI/(N-1); //Sets up the difference between each

```

```

//x values, all equall spaced

ofstream *file = NULL;    //used create the file the data
                           //is stored in.

ofstream *file1 = NULL;
ofstream *file2 = NULL;
ofstream *file3 = NULL;

vector<double*> data_store;    //a vector to store all pigmentation
                              //patterns for all time iterations

vector<double*> data_store_1;  //a vector to store all the
                              //linearization values

vector<double> dist;          //stores the infinity norm for
                              //each iteration

vector<double> norm_dist;

bool testflag = true;

enum comp { real, imaginary }; //this is used to help me keep track of the
                               //multiplication of complex numbers.

int main(int argc, char *argv[])
{
    string yes_no = "";
    srand((int)time(NULL));
    initializationVariables();
    do
    {
        cout << "here again" << endl;
        getInput();
        computeEquilibrium();
    }
}

```

```

    patternGen();

    cout << "Continue?[y/n]: ";

    cin >> yes_no;

    cout << "here" << endl;

    while( yes_no[0] != 'y' && yes_no[0] != 'n' )
    {
        cout << "Invalid Entry" << endl

            << "Continue?[y/n]:";

        cin >> yes_no;
    }

}while( yes_no[0] != 'n' );

cleanUp();

return 0;
}

/*The main functions******/

void initializationVariables()
{
    cerr << "Variable Initialization: ";

    L0 = (double*)fftw_malloc(N*sizeof(double));

    L0t = (double*)fftw_malloc(N*sizeof(double));

    L0t_p_1 = (double*)fftw_malloc(N*sizeof(double));

    cerr << "*";

    L0t_p_2 = (double*)fftw_malloc(N*sizeof(double));

    in_Ut = (double*)fftw_malloc(N*sizeof(double));

    out_Ut = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));

    cerr << "*";

```

```

in_wE = (double*)fftw_malloc(N*sizeof(double));
in_wI = (double*)fftw_malloc(N*sizeof(double));
out_wE = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
cerr << "*";
out_wI = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
in_Pt = (double*)fftw_malloc(N*sizeof(double));
out_Pt = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
cerr << "*";
in_SwE = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
out_SwE = (double*)fftw_malloc(N*sizeof(double));
in_SwI = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
cerr << "*";
out_SwI = (double*)fftw_malloc(N*sizeof(double));
in_SwE1 = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
out_SwE1 = (double*)fftw_malloc(N*sizeof(double));
cerr << "*";
in_SwI1 = (fftw_complex*)fftw_malloc(N*sizeof(fftw_complex));
out_SwI1 = (double*)fftw_malloc(N*sizeof(double));
Ut = (double*)fftw_malloc(N*sizeof(double));
cerr << "*";
Ut_p_1 = (double*)fftw_malloc(N*sizeof(double));
Ut_p_2 = (double*)fftw_malloc(N*sizeof(double));
Pt = (double*)fftw_malloc(N*sizeof(double));
cerr << "*";
Pt_p_1 = (double*)fftw_malloc(N*sizeof(double));
Pt_p_2 = (double*)fftw_malloc(N*sizeof(double));
fft_wE = fftw_plan_dft_r2c_1d(N, in_wE, out_wE, FFTW_EXHAUSTIVE);

```

```

    cerr << "*";

    fft_wI = fftw_plan_dft_r2c_1d(N, in_wI, out_wI, FFTW_EXHAUSTIVE);
    fft_Pt = fftw_plan_dft_r2c_1d(N, in_Pt, out_Pt, FFTW_EXHAUSTIVE);
    fft_Ut = fftw_plan_dft_r2c_1d(N, in_Ut, out_Ut, FFTW_EXHAUSTIVE);

    cerr << "*";

    ifft_SwE = fftw_plan_dft_c2r_1d(N, in_SwE, out_SwE, FFTW_EXHAUSTIVE);
    ifft_SwI = fftw_plan_dft_c2r_1d(N, in_SwI, out_SwI, FFTW_EXHAUSTIVE);
    ifft_SwE1 = fftw_plan_dft_c2r_1d(N, in_SwE1, out_SwE1, FFTW_EXHAUSTIVE);

    cerr << "*";

    ifft_SwI1 = fftw_plan_dft_c2r_1d(N, in_SwI1, out_SwI1, FFTW_EXHAUSTIVE);

    precision = 0.5*pow(10, -1*accuracy);

    cerr << endl << "Initialization Complete" << endl;

    return;
}

void getInput()
{
    string temp;

    cerr << "Enter theta_E[" << theta_E << "]: ";

    if( testflag )
        testflag = false;
    else
        cin >> ws;

    getline(cin, temp, '\n');

    if( temp != "" )
    {
        theta_E = atof( temp.c_str() );
    }
}

```

```

cerr << "Enter theta_I[" << theta_I << "]: ";
getline(cin, temp, '\n');
if( temp != "" )
{
    theta_I = atof( temp.c_str() );
}
cerr << "Enter alpha_E[" << alpha_E << "]: ";
getline(cin, temp, '\n');
if( temp != "" )
{
    alpha_E = atof( temp.c_str() );
}
cerr << "Enter alpha_I[" << alpha_I << "]: ";
getline( cin, temp, '\n');
if( temp != "" )
{
    alpha_I = atof( temp.c_str() );
}
cerr << "Enter sigma_E[" << sigma_E << "]: ";
getline( cin, temp, '\n');
if( temp != "" )
{
    sigma_E = atof( temp.c_str() );
}
cerr << "Enter sigma_I[" << sigma_I << "]: ";
getline( cin, temp, '\n');
if( temp != "" )

```

```

{
    sigma_I = atof( temp.c_str() );
}

cerr << "Enter gamma[" << gamma1 << "]: ";
getline( cin, temp, '\n' );
if( temp != "" )
{
    gamma1 = atof( temp.c_str() );
}

cerr << "Enter delta[" << delta << "]: ";
getline( cin, temp, '\n' );
if( temp != "" )
{
    delta = atof( temp.c_str() );
}

cerr << "Enter nu_E[" << nu_E << "]: ";
getline( cin, temp, '\n' );
if( temp != "" )
{
    nu_E = atof( temp.c_str() );
}

cerr << "Enter nu_I[" << nu_I << "]: ";
getline( cin, temp, '\n' );
if( temp != "" )
{
    nu_I = atof( temp.c_str() );
}

```



```

}

void cleanUp()
{
    fftw_free(L0);
    fftw_free(L0t);
    fftw_free(L0t_p_1);
    fftw_free(L0t_p_2);
    fftw_free(in_wI);
    fftw_free(out_wE);
    fftw_free(out_wI);
    fftw_free(in_Pt);
    fftw_free(out_Pt);
    fftw_free(in_SwE);
    fftw_free(out_SwE);
    fftw_free(in_SwI);
    fftw_free(out_SwI);
    fftw_free(in_SwE1);
    fftw_free(out_SwE1);
    fftw_free(in_SwI1);
    fftw_free(out_SwI1);
    fftw_destroy_plan(fft_wE);
    fftw_destroy_plan(fft_wI);
    fftw_destroy_plan(fft_Pt);
    fftw_destroy_plan(fft_Ut);
    fftw_destroy_plan(ifft_SwE);
    fftw_destroy_plan(ifft_SwI);

```

```

    fftw_destroy_plan(iff_t_SwE1);

    fftw_destroy_plan(iff_t_SwI1);

} //end of cleanUp()

void patternGen()
{
    cerr << "Calculations Begin: ";

    q_E = alpha_E*4/(93*sigma_E);

    q_I = alpha_I*4/(93*sigma_I);

    Pt = (double*)fftw_malloc(N*sizeof(double));

    Pt_p_1 = (double*)fftw_malloc(N*sizeof(double));

    cerr << "*";

    Pt_p_2 = (double*)fftw_malloc(N*sizeof(double));

    if( file != NULL )
    {
        delete file;

        delete file1;

        delete file2;

        delete file3;
    }

    file = new ofstream();

    file1 = new ofstream();

    file2 = new ofstream();

    file3 = new ofstream();

    cerr << "*";

    file->open("/media/SEA__DISK/outfile.dat");

    file1->open("/media/SEA__DISK/outfile1.dat");

    file2->open("/media/SEA__DISK/outfile2.dat");

```

```

file3->open("/media/SEA__DISK/outfile3.dat");

(*file) << "%theta_E: " << theta_E << " theta_I: " << theta_I << endl

                                //writes variable values to file

    << "%alpha_E: " << alpha_E << " alpha_I: " << alpha_I << endl
    << "%sigma_E: " << sigma_E << " sigma_I: " << sigma_I << endl
    << "%gamma: " << gamma1 << " delta: " << delta << endl
    << "%nu_E: " << nu_E << " nu_I: " << nu_I << endl
    << "%equilibrium: " << equilibrium << endl;

(*file1) << "%theta_E: " << theta_E << " theta_I: " << theta_I << endl

    << "%alpha_E: " << alpha_E << " alpha_I: " << alpha_I << endl
    << "%sigma_E: " << sigma_E << " sigma_E: " << sigma_I << endl
    << "%gamma: " << gamma1 << " delta: " << delta << endl
    << "%nu_E: " << nu_E << " nu_I: " << nu_I << endl
    << "%equilibrium: " << equilibrium << endl;

(*file2) << "%theta_E: " << theta_E << " theta_I: " << theta_I << endl

                                //writes variable values to file

    << "%alpha_E: " << alpha_E << " alpha_I: " << alpha_I << endl
    << "%sigma_E: " << sigma_E << " sigma_I: " << sigma_I << endl
    << "%gamma: " << gamma1 << " delta: " << delta << endl
    << "%nu_E: " << nu_E << " nu_I: " << nu_I << endl
    << "%equilibrium: " << equilibrium << endl;

for( int i = 0; i < N; i++ )
{
                                //initializes the variables
    x_coor[i] = diff*i;

    Ut[i] = ((double)(rand()%1000-500.0)/500000.0);

                                //chooses random points

    Ut_p_1[i] = ((double)(rand()%1000-500.0)/500000.0);

```

```

                                                                    //near the equilibrium

    Pt[i] = equilibrium+Ut[i];                                //chooses random points

    Pt_p_1[i] = equilibrium+Ut_p_1[i];                        //near zero

    in_wE[i] = wE(x_coor[i]);

    in_wI[i] = wI(x_coor[i]);

}

cerr << "*";

data_store.push_back( Pt );                                //stores the first

data_store.push_back( Pt_p_1 );                            //and second iteration

data_store_1.push_back( Ut );                                //of Pt, Ut and

data_store_1.push_back( Ut_p_1 );                            //their difference

dist.push_back( infinity_norm( Pt, Ut ) );

norm_dist.push_back(rel_norm( Pt, Ut, dist[0] ) );

cerr << "*";

dist.push_back( infinity_norm( Pt_p_1, Ut_p_1 ) );

norm_dist.push_back(rel_norm( Pt_p_1, Ut_p_1, dist[1] ) );

arrayCpy( in_Pt, Pt );

arrayCpy( in_Ut, Ut );

fftw_execute(fft_Pt);

cerr << "*";

fftw_execute(fft_wE);

fftw_execute(fft_wI);

computeL0();

ofstream test;

cerr << "*";

computeConvolution( ifft_SwE, out_Pt, out_wE, in_SwE );

computeConvolution( ifft_SwI, out_Pt, out_wI, in_SwI );

```

```

arrayCpy( L0t, L0 );

cerr << "*";

for( int i = 2; i < max_time; i++ )
{
    arrayCpy( in_Ut, Ut_p_1 );
    arrayCpy( in_Pt, Pt_p_1 );
    fftw_execute(fft_Pt);
    computeL0();
    arrayCpy( L0t_p_1, L0 );
    arrayCpy( temp_SwE, out_SwE );
    arrayCpy( temp_SwI, out_SwI );
    computeConvolution( ifft_SwE, out_Pt, out_wE, in_SwE );
    computeConvolution( ifft_SwI, out_Pt, out_wI, in_SwI );
    for( int j = 0; j < N; j++ )
    {
        Pt_p_2[j] = S(out_SwE[j]*2*PI/(N*N),out_SwI[j]*2*PI/(N*N)) -
            gamma1*Pt[j] -
            delta*( S(temp_SwE[j]*2*PI/(N*N),temp_SwI[j]*2*PI/(N*N)) -
                Pt_p_1[j]);
        Ut_p_2[j] = L0t_p_1[j] + delta*Ut_p_1[j] - delta*L0t[j] -
            gamma1*Ut[j];
    }
    data_store_1.push_back(Ut_p_2);
    data_store.push_back(Pt_p_2);
    dist.push_back(infinity_norm(Pt_p_2, Ut_p_2));
    norm_dist.push_back(rel_norm(Pt_p_2, Ut_p_2, dist[i]));
    Ut = Ut_p_1;
}

```

```

    Pt = Pt_p_1;

    Ut_p_1 = Ut_p_2;

    Pt_p_1 = Pt_p_2;

    Ut_p_2 = (double*)fftw_malloc(N*sizeof(double));

    Pt_p_2 = (double*)fftw_malloc(N*sizeof(double));

    arrayCpy(L0t, L0t_p_1);

}

cerr << "*";

for( int i = 0; i < max_time; i++ )
{
    (*file) << dist[i] << " ";

    (*file3) << norm_dist[i] << " ";

}

(*file) << endl;

(*file3) << endl;

for( int i = 0; i < 2000; i++ )
{
    for( int j = 0; j < N; j++ )

        (*file1) << data_store[i][j] << " ";

    (*file1) << endl;

}

for( int i = 0; i < 2000; i++ )
{
    for( int j = 0; j < N; j++ )

        (*file2) << data_store_1[i][j] << " ";

    (*file2) << endl;

}

```

```

for( int i = 0; i < max_time; i++ )
{
    fftw_free(data_store[i]);

    fftw_free(data_store_1[i]);
}

cerr << "*";

data_store.clear();

data_store_1.clear();

dist.clear();

norm_dist.clear();

file->close();

file1->close();

file2->close();

file3->close();

delete file1;

delete file2;

delete file;

delete file3;

file3 = NULL;

file1 = NULL;

file2 = NULL;

file = NULL;

cerr << endl << "Calculations Complete" << endl;
}

/*End of the main functions*****

/*These functions are used to calculate S[Pt(x)]*****

```

```

double wE( double x )
{
    if( x > sigma_E && x < (2*PI-sigma_E) )
        return 0;
    else
    {
        if( x < sigma_E )
            return (q_E * ( ( 1 << p) - pow(1 - cos(PI*x/sigma_E), p )));
        else
        {
            return (q_E * ( ( 1 << p) -
                pow(1 - cos(PI*(2*PI-x)/sigma_E), p )));
        }
    }
}

double wI( double x )
{
    if( x > sigma_I && x < (2*PI-sigma_I) )
        return 0;
    else
    {
        if( x < sigma_I )
            return (q_I * ( ( 1 << p ) - pow(1 - cos(PI*x/sigma_I), p )));
        else
        {
            return (q_I * ( ( 1 << p ) -
                pow(1 - cos(PI*(2*PI-x)/sigma_I), p )));
        }
    }
}

```



```

    }
}

double SE( double x )
{
    return (1/(1+exp(-1*nu_E*(x - theta_E))));
}

double SI( double x )
{
    return (1/(1+exp(-1*nu_I*(x - theta_I))));
}

double S( double x_E, double x_I )
{
    return (SE(x_E)-SI(x_I));
}

/*End of the calculations used to S[Pt(x)]*****/

/*This sections is used to calculate the homogenous equilibrium *****/

double newtonMethodP0( double x )
{
    double temp;

    temp = x - equationNewton( x )/equationNewtonPrime( x );

    return temp;
}

```

```

double equationNewton( double x )
{
    double temp;

    temp = x - (1-delta)/(1+gamma1-delta)*(SE(alpha_E*x) - SI(alpha_I*x));

    return temp;
}

double equationNewtonPrime( double x )
{
    double temp;

    temp = 1 - (1-delta)/(1+gamma1-delta)*(alpha_E*SprimeE(alpha_E*x)
    - alpha_I*SprimeI(alpha_I*x));

    return temp;
}

void computeEquilibrium()
{
    double xt = 0, xt_p_1 = 0;

    do
    {
        xt = xt_p_1;

        xt_p_1 = newtonMethodP0(xt);

    }while( fabs(xt - xt_p_1) > precision );

    equilibrium = xt_p_1;

    cerr << "Equilibrium: " << equilibrium << endl;
}

/*End of the section for computing the homogenous equilibrium*****/

```

```
/*These functions are only used in the linearization calculations*****/
```

```
double SprimeE( double x )
```

```
{
    return (nu_E*exp(-1*nu_E*(x - theta_E))/pow(1+exp(-1*nu_E*(x -
    theta_E)),2));
}
```

```
double SprimeI( double x )
```

```
{
    return (nu_I*exp(-1*nu_I*(x - theta_I))/pow(1+exp(-1*nu_I*(x -
    theta_I)),2));
}
```

```
void computeL0()
```

```
{
    fftw_execute(fft_Ut);
    computeConvolution( ifft_SwE1, out_Ut, out_wE, in_SwE1 );
    computeConvolution( ifft_SwI1, out_Ut, out_wI, in_SwI1 );
    for( int i = 0; i < N; i++ )
        L0[i] = SprimeE(alpha_E*equilibrium)*out_SwE1[i]*2*PI/(N*N) -
        SprimeI(alpha_I*equilibrium)*out_SwI1[i]*2*PI/(N*N);
    return;
}
```

```
/*End of the functions specifically used in the linearization calculations**/
```

```
/*These are the auxillary functions used to make the code more modular*****/
```

```

void arrayCpy( double a[], double b[] )
{
    for( int i = 0; i < N; i++ )
        a[i] = b[i];
}

void computeConvolution(fftw_plan &plan, fftw_complex array1[],
                        fftw_complex array2[], fftw_complex array3[])
{
    for( int i = 0; i < N; i++ )
    {
        array3[i][(int)real] = array1[i][(int)real]*array2[i][(int)real] -
        array1[i][(int)imaginary]*array2[i][(int)imaginary];
        array3[i][(int)imaginary] = array1[i][(int)real]*
        array2[i][(int)imaginary] +
        array1[i][(int)imaginary]*array2[i][(int)real];
    }
    fftw_execute(plan);
}

double infinity_norm(double *pt, double *ut)
{
    double maximum = 0;
    for( int i = 0; i < N; i++ )
        maximum = max(fabs((pt[i]-equilibrium)-ut[i]), maximum) ;
    return maximum;
}

```

```
double rel_norm(double *pt, double *ut, double numerator)
{
    double maximum = 0;

    for( int i = 0; i < N; i++ )
        maximum = max(fabs(ut[i]), maximum);

    return numerator/maximum;
}

/*****
```

## A.2 Maple Code

```
> restart;
> p:=4;thetae:=4;thetai:=3.3;alpe:=15; alpi:=13;sigmaE:=1.3;
    sigmaI:=.4;gam:=.2; delta:=.1; vue:=.4; vui:=.4;
> plot([sqrt(gam)+del, 1+gam/(1+del), 1+gam/(1-del),
    (1-gam)/del],del=0..1,y=0..2, legend=["1","2","3","4"]);
> sqrt(gam)+delta; 1+gam/(1+delta); 1+gam/(1-delta); (1-gam)/delta;
> f1:=2^p-(1-cos(Pi*x/sigmaE))^p;
> f2:=2^p-(1-cos(Pi*x/sigmaI))^p;
> qE:=alpe/int(f1,x=-sigmaE..sigmaE);
> qI:=alpi/int(f2,x=-sigmaI..sigmaI);
> we:=f1*qE; wi:=f2*qI;
> int(we,x=-sigmaE..sigmaE);int(wi,x=-sigmaI..sigmaI);
> plot({we,wi},x=-sigmaE..sigmaE);
> evalf(subs(x=sigmaE/2,we));evalf(subs(x=sigmaI/2,wi));
```

```

> plot(wi,x=-sigmaI..sigmaI);
> fte:=(int(we*exp(I*k*x),x=-sigmaE..sigmaE)):
> fti:=(int(wi*exp(I*k*x),x=-sigmaI..sigmaI)):
> evalf(subs(k=1.00001,fte));
> plot({fte,fti},k=0..100);
> se:=x->1/(1+exp(-vue*(x-thetae)));
> si:=x->1/(1+exp(-vui*(x-thetai)));
> plot({p1,(1-delta)/(1+gam-delta)*(se(p1*alpe)-
    si(p1*alpi))},p1=-1..1);
> p00:=fsolve(((1-delta)/(1+gam-delta)*(se(p0*alpe)-
    si(p0*alpi))=p0,p0=0);
> seprime:=evalf(subs(x=alpe*p00,diff(se(x),x)));
> siprime:=evalf(subs(x=alpi*p00,diff(si(x),x)));
> lstar:=(seprime*fte-siprime*fti):
> plot({lstar},k=0..150);
> lmin:=limit(lstar,k=0);
> a:=-Re(lstar)-delta:
> b:=delta*Re(lstar)+gam:
> sominus:=(-a-sqrt(a^2-4*b))/2:
> soplus:=(-a+sqrt(a^2-4*b))/2:
> plot({1,abs(soplus),abs(sominus)},k=0..300);
> with(plots):
> plot1:=plot(x-1,x=0..2,y=-1.5..1.5):
> plot2:=plot(-x-1,x=-2..0):
> plot3:=plot(1,x=-2..2):

```

```

> plot4:=plot([a,b,k=0..10],color=blue):
> plot5:=plot([a,b,k=10..100],color=green):
> plot6:=plot(.25*k^2,k=-2..2):
> display({plot1,plot2,plot3,plot4,plot5,plot6});

```

## A.3 Matlab Code

### A.3.1 Nonlinear Pattern Grapher

```

x = linspace(0,2*pi,4096);
y = linspace(900,1000,101);
pcolor(x,y,outfile1(900:1000,:));
shading interp;
%caxis([0.038, .039]);
colorbar;
set(gca,'fontsize',15);
xlabel('Space','fontsize',16);
ylabel('Time','fontsize',16);
title('Nonlinear Bifurcation Thru Complex', 'fontsize', 16);

```

### A.3.2 Linear Pattern Grapher

```

x = linspace(0,2*pi,4096);
y = linspace(500,520,21);
pcolor(x,y,outfile2(500:520,:));
shading interp;
%caxis([-1 1]);

```

```
colorbar;  
set(gca,'fontsize',15);  
xlabel('Space','fontsize',16);  
ylabel('Time','fontsize',16);  
title('Linear Bifurcation Thru Complex', 'fontsize', 16);
```